



# Entropy in nearby galaxy clusters



Meredith Reitz

3/9

*XMM-Newton observations of three poor clusters: Similarity in dark matter and entropy profiles down to low mass*

[Pratt, G. and Arnaud, M. , 2005]

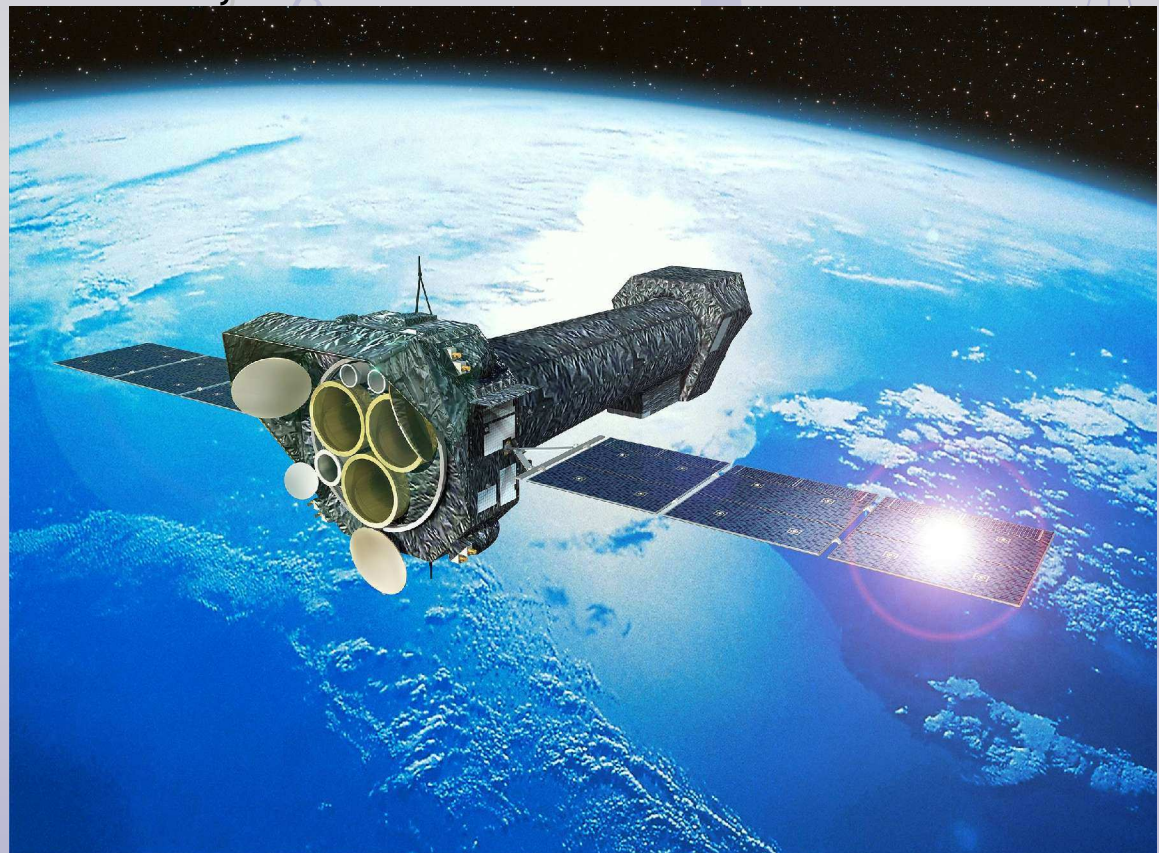
**&**

*Structure and scaling of the entropy in nearby galaxy clusters*

[Pratt, G., Arnaud, M., & Pointecouteau, E. , 2006]

# XMM-Newton

- NASA-ESA
- Launched 12 / 1999, still active
- Observes  
x-rays of  
0.2-12 keV





*First paper:*

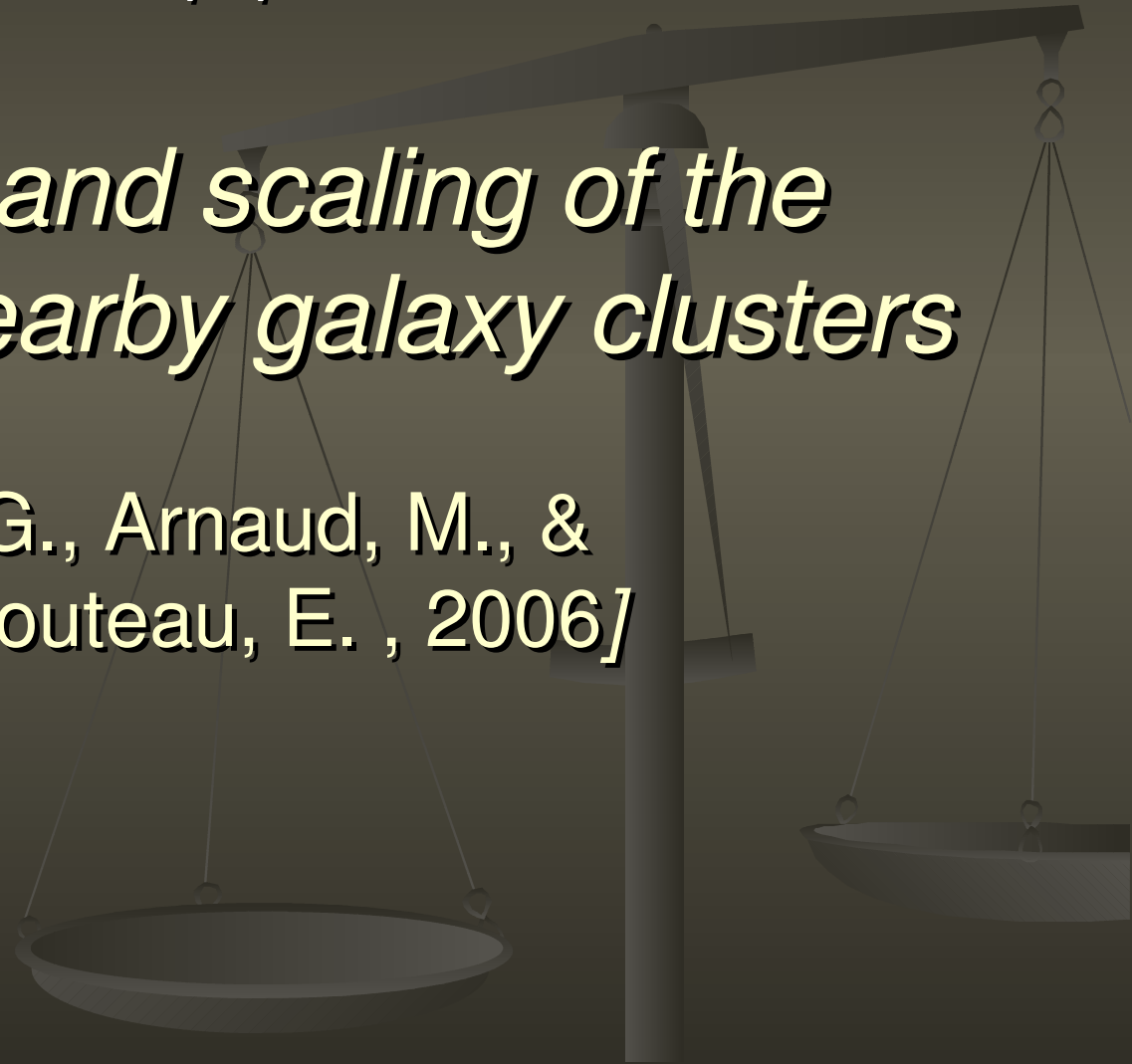
*XMM-Newton observations of three  
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*Second paper:*

# *Structure and scaling of the entropy in nearby galaxy clusters*

[Pratt, G., Arnaud, M., &  
Pointecouteau, E. , 2006]



# Motivation

- Better understanding of cluster formation / evolution
- See to what degree non-gravitational processes are significant
- Mass profile  $M(r)$ 
  - Information about gravitational collapse
- Entropy is generated in shocks as gas is drawn into the potential well of the cluster
- Entropy profile  $S(r)$ 
  - ICM accretion, thermodynamic history
  - Non-gravitational processes

# These x-rays

- Intracluster medium (ICM) – hot gas
- Two quantities define x-ray properties
  - Entropy profile of the gas,  $S(r)$
  - Shape of the gravitational potential well;  $M(r)$
- Low-mass clusters
  - Non-gravitational, gravitational comparable

# This paper

- Sample: A1991, A2717, MKW9
  - Three low-mass, cool clusters
  - $0.04 \leq z \leq 0.06$
  - $kT = 2.65, 2.53, 2.58$  keV
  - Combine with previous results for cool A1983 ( $kT = 2.2$  keV) and hot A1413 ( $kT = 6.5$  keV)
- Assumptions
  - $\Lambda$ CDMH70:  $H_0 = 70$  km s<sup>-1</sup> Mpc<sup>-1</sup>,  $\Omega_m = 0.3$ ,  $\Omega_\Lambda = 0.7$
  - Some SCDMH50:  $H_0 = 50$  km s<sup>-1</sup> Mpc<sup>-1</sup>,  $\Omega_m = 1.0$ ,  $\Omega_\Lambda = 0.0$

# Sample of clusters

- 10 systems
  - Temperatures:  
 $kT = 2 \text{ keV} \dots 8.5 \text{ keV}$
  - Redshifts:  $0.03 \leq z \leq 0.15$
- $\Lambda$ CDMH70

Cluster	$z$	$T$ (keV)
A1983	0.0442	$2.18 \pm 0.09$
A2717	0.0498	$2.56 \pm 0.06$
MKW9	0.0382	$2.43 \pm 0.24$
A1991	0.0586	$2.71 \pm 0.07$
A2597	0.0852	$3.67 \pm 0.09$
A1068	0.1375	$4.67 \pm 0.11$
A1413	0.1427	$6.62 \pm 0.14$
A478	0.0881	$7.05 \pm 0.12$
PKS0745	0.1028	$7.97 \pm 0.28$
A2204	0.1523	$8.26 \pm 0.22$

# Overview

- Expectations from models of cluster formation, discrepancies
  - \*Mass and entropy
- Surface brightness profiles
  - Emission measure profiles
  - Gas density profiles
- Hardness ratio images
  - Temperature structure – dynamical state
- Annular spectral analysis & surface brightness profiles
  - Abundance profiles
  - Temperature profiles
    - Correct for projection, PSF effects (gas density profiles)
- Mass and entropy profiles
  - From gas density, temperature profiles
  - And scaled



# Overview

- Larger sample size, - scaling parameters
- Entropy-temperature relation
  - At different fractions of virial radius
    - $R_{200}$ : radius within which density is 200 times  $\rho_c(z)$ :  
$$\rho_c(z) = 3h(z)^2 H_0 / 8\pi G \quad h^2(z) = \Omega_m(1+z)^3 + \Omega_\Lambda$$
- Entropy-mass relation
- Entropy profiles,  $S(r)$ 
  - Scaled with best-fitting  $S - T$ ,  $S - M$  relations

# Self-similar models

- Predict universal shape for dark matter distribution from  $\sim 0.01 r_{200}$  to  $\sim 0.7 r_{200}$ , high and low mass clusters
- Predict central cusp
- X-ray observations can confirm
- Cluster formation governed solely by gravitational processes
- Imply that properties  $S$ ,  $L_x$  would scale with powers of  $T$ ,  $M$
- Scaling is somewhat inconsistent with observations

# Self-similar models

- Entropy (pseudoentropy) defined as

$$S = kT/n_e^{2/3}$$

scales theoretically with  
temperature as

$$S \propto h(z)^{-4/3} T$$

- But scales empirically closer to

$$S \propto T^{0.65}$$

- {second paper:  $0.64 \pm 0.11$ }

- With radius, in model, scales as

$$S \propto r^{1.1}$$

- But scales empirically like  
(with  $T^{0.65}$  scaling)

$$S \propto r^{0.94 \pm 0.14}$$

- $\{1.08 \pm 0.04\}$

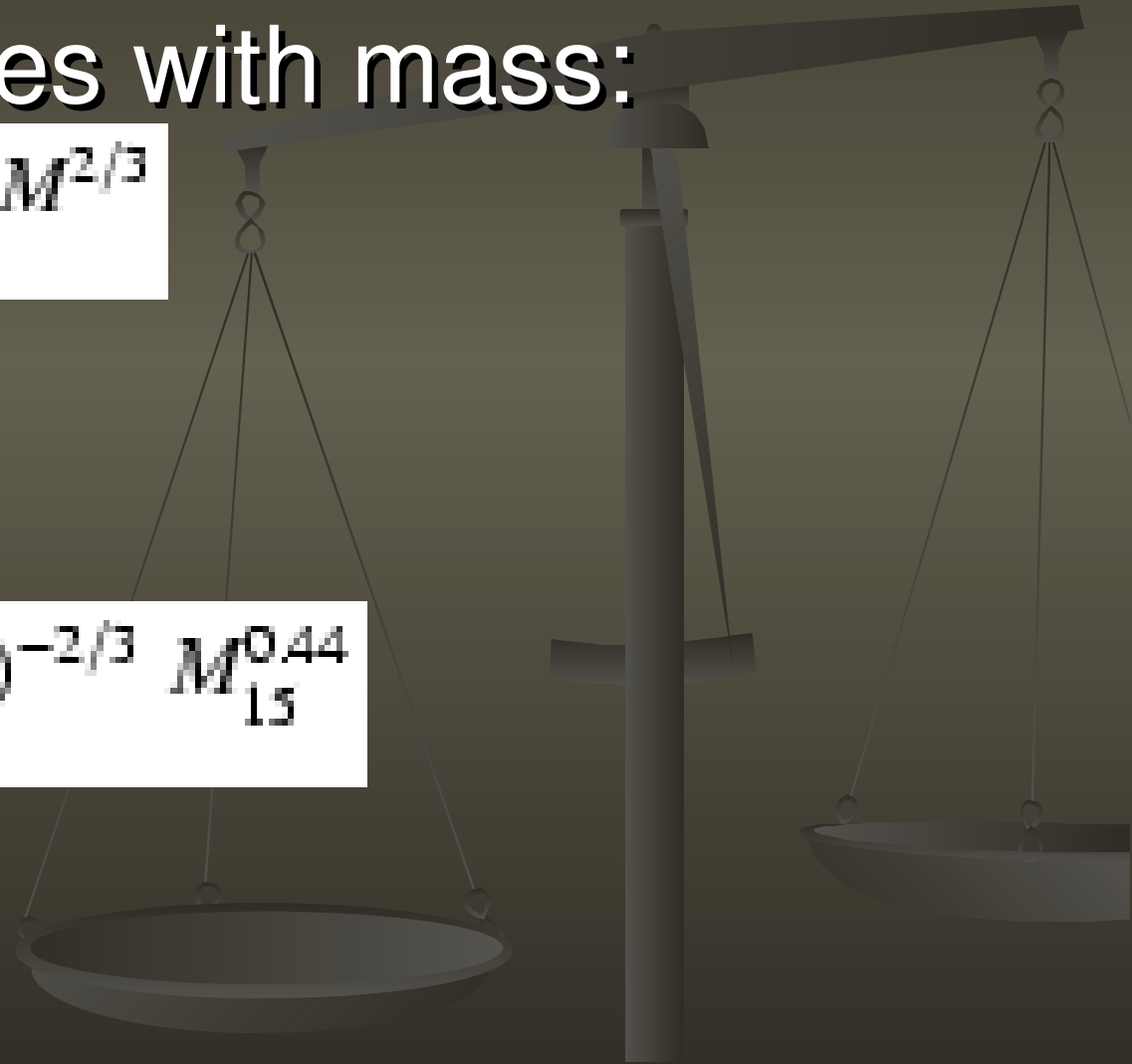
# Self-similar models

- Entropy scales with mass:

$$S \propto h(z)^{-2/3} M^{2/3}$$

- Find:

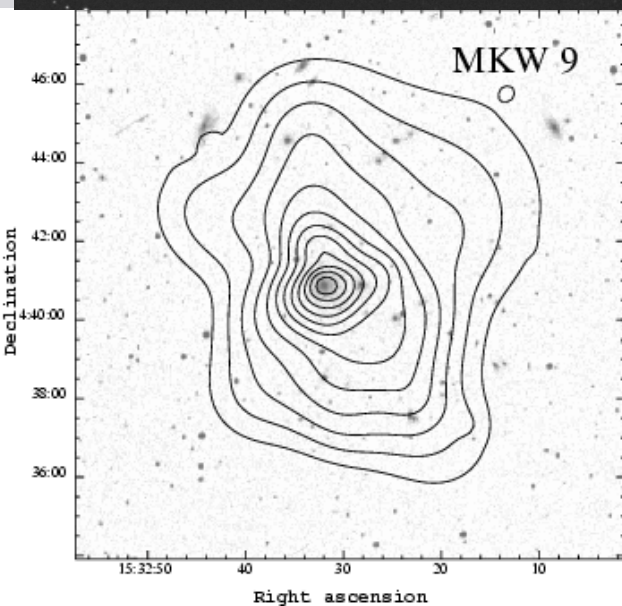
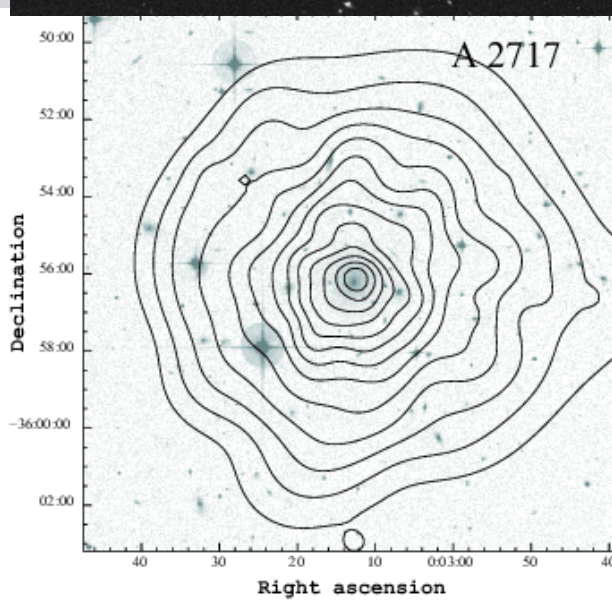
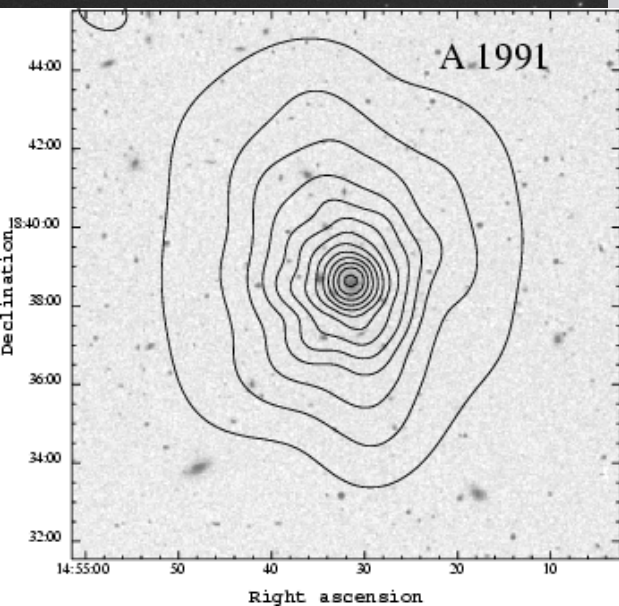
$$S_{2500} \propto h(z)^{-2/3} M_{15}^{0.44}$$



# Possible non-gravitational processes

- Preheating of gas before accretion
  - Early supernovae
    - 'seems too localized to have a significant effect in smoothing the accreting gas'
  - Early AGN activity
  - Only is unlikely for lack of observed isentropic cores
- Internal heating after accretion
- Cooling
  - Predicts higher stellar mass fraction than observed
- Likely
  - Interplay between cooling and feedback, combined with some preheating

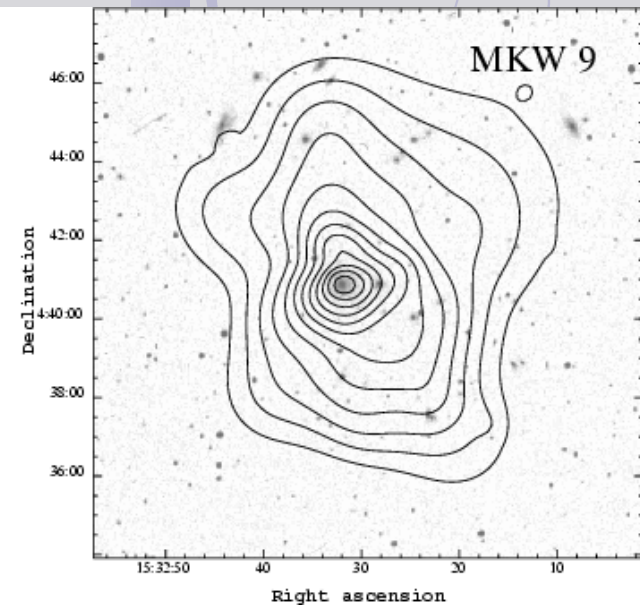
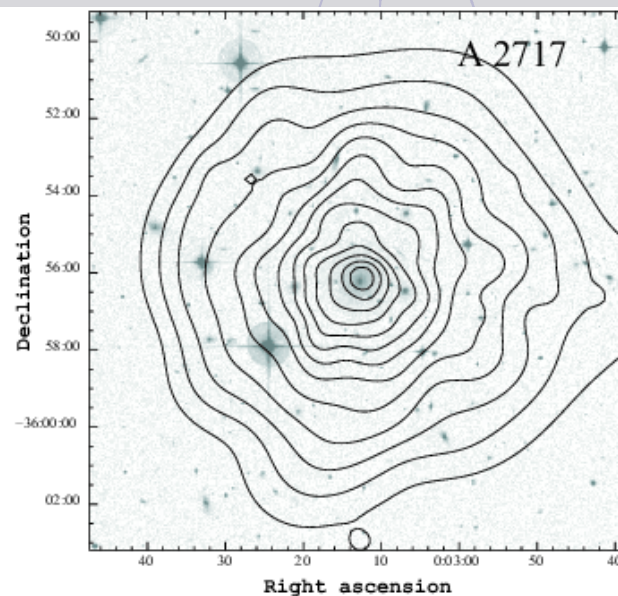
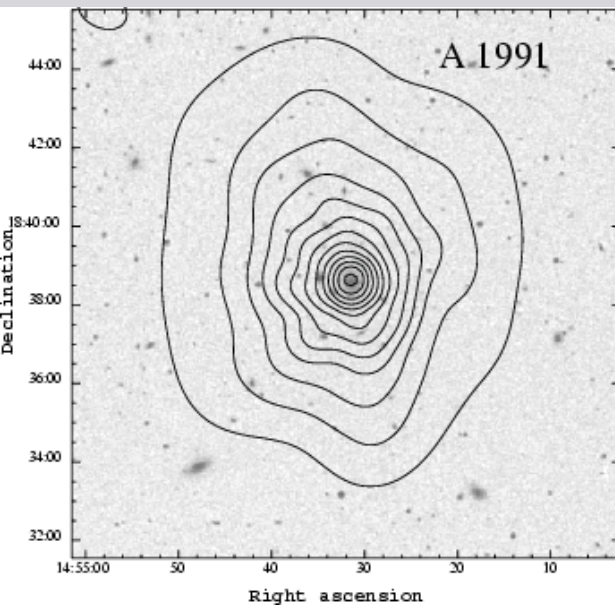
# X-ray morphologies





# X-ray morphologies

- A1991 and A2717: symmetric x-ray isophotes → relatively relaxed
- MKW9 asymmetrical
  - A2717 centered on central galaxy ACO 2717 BCG
  - A1991 somewhat off-center from central galaxy ACO 1991A
  - MKW9 centered on central galaxy UGC 9886

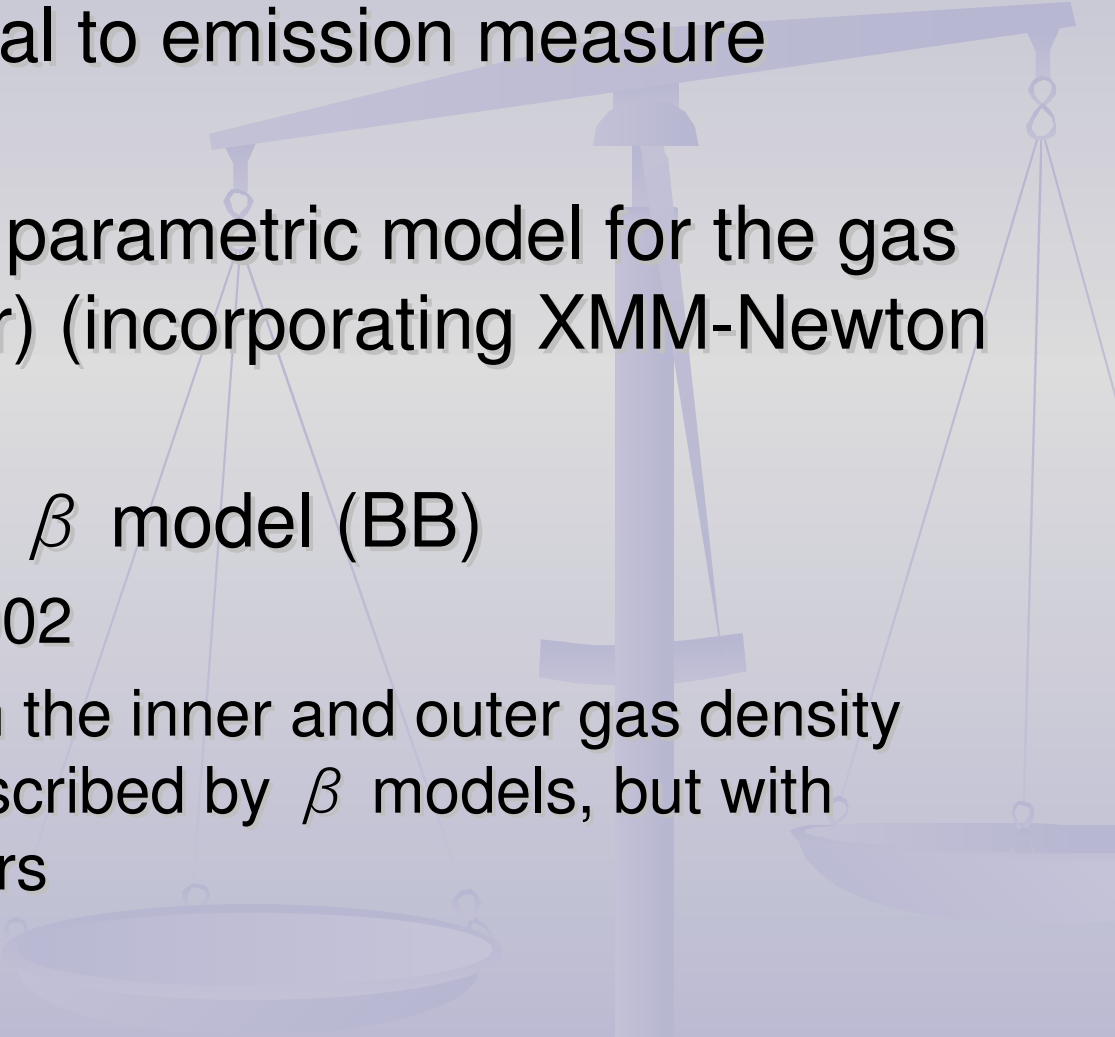




# Surface brightness profiles

- $\Lambda$  – emissivity of the hot gas – depends sensitively on gas abundance / temperature
- Data reduction
  - Masking of point sources
  - Looking at 0.3 – 3.0 keV bands
  - Corrected for emissivity variations - temperature / abundance profiles fitted to functional forms
  - $\Lambda(\text{radius})$  estimated with a MEKAL model (a thermal equilibrium plasma emission model), normalized to its value at large radius

# Surface brightness profiles

- Directly proportional to emission measure profile,  $EM(r)$
  - Can be fitted from parametric model for the gas density profile,  $n_e(r)$  (incorporating XMM-Newton PSF)
  - Double isothermal  $\beta$  model (BB)
    - Pratt & Arnaud, 2002
    - Assumes that both the inner and outer gas density profiles can be described by  $\beta$  models, but with different parameters
- 

# Gas density profile - double isothermal $\beta$ model

- $n_H(r)$  = the gas density radial profile
- $R_{\text{cut}}$  = free parameter

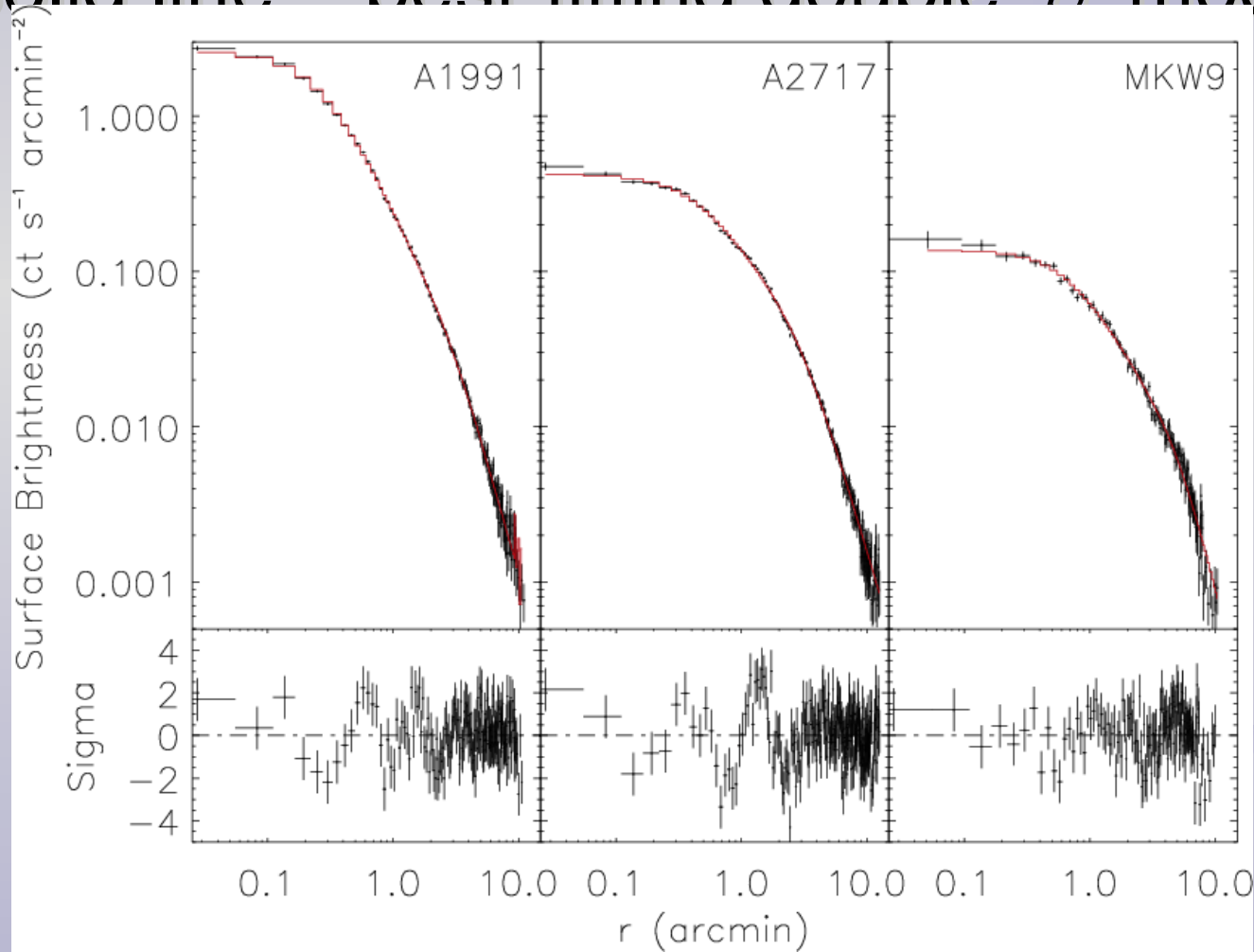
$$\begin{aligned} r < R_{\text{cut}} \quad n_H(r) &= n_{H,0} \left[ 1 + \left( \frac{r}{r_{c,n}} \right)^2 \right]^{-\frac{3\beta_{in}}{2}} \\ r > R_{\text{cut}} \quad n_H(r) &= N \left[ 1 + \left( \frac{r}{r_c} \right)^2 \right]^{-\frac{3\beta}{2}} \end{aligned}$$

$$\beta_{in} = \beta \frac{1 + \left( \frac{r_{c,n}}{R_{\text{cut}}} \right)^2}{1 + \left( \frac{r_c}{R_{\text{cut}}} \right)^2}$$

$$N = n_{H,0} \frac{\left[ 1 + \left( \frac{R_{\text{cut}}}{r_{c,n}} \right)^2 \right]^{-\frac{3\beta_{in}}{2}}}{\left[ 1 + \left( \frac{R_{\text{cut}}}{r_c} \right)^2 \right]^{-\frac{3\beta}{2}}}$$

# Surface brightness profiles

Solid line – best-fitting double  $\beta$  model

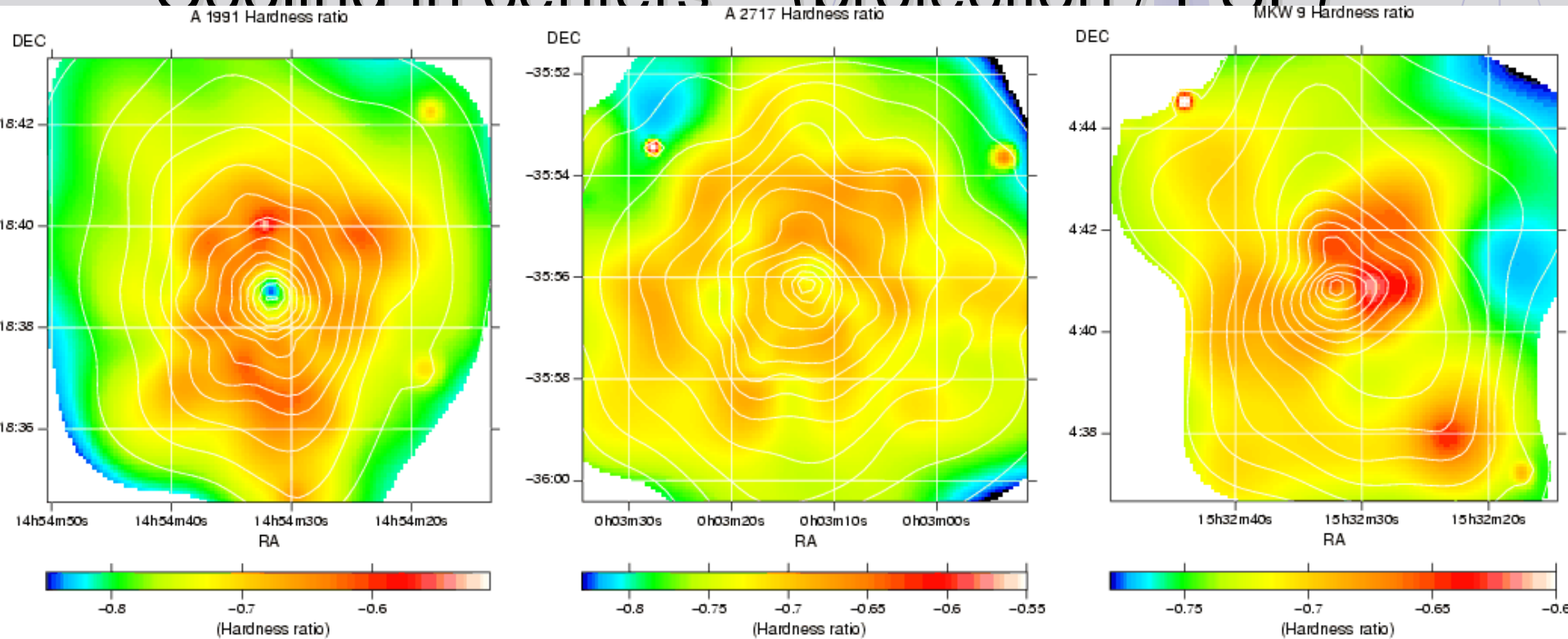


# Temperature distributions – hardness ratio images

- Hardness ratio: ratio of counts in different wave bands – a measure of the spectral slope of a source – an indirect measure of temperature
- Source / background images subtracted, smoothed images with smoothing scale of  $2.5 \sigma$  -  $4 \sigma$
- *Not* corrected for difference between local cluster backgrounds and blank-sky backgrounds at low energies
  - HR values cannot be converted directly, reliably into temperatures
  - HR decline toward outer regions an artifact
- Valuable for understanding the temperature structure

# Temperature distributions – hardness ratio images

- Again, A1991 and A2717 symmetrical
- MKW9 asymmetrical – not entirely relaxed
- Cooling in centers – (projection / PSF)

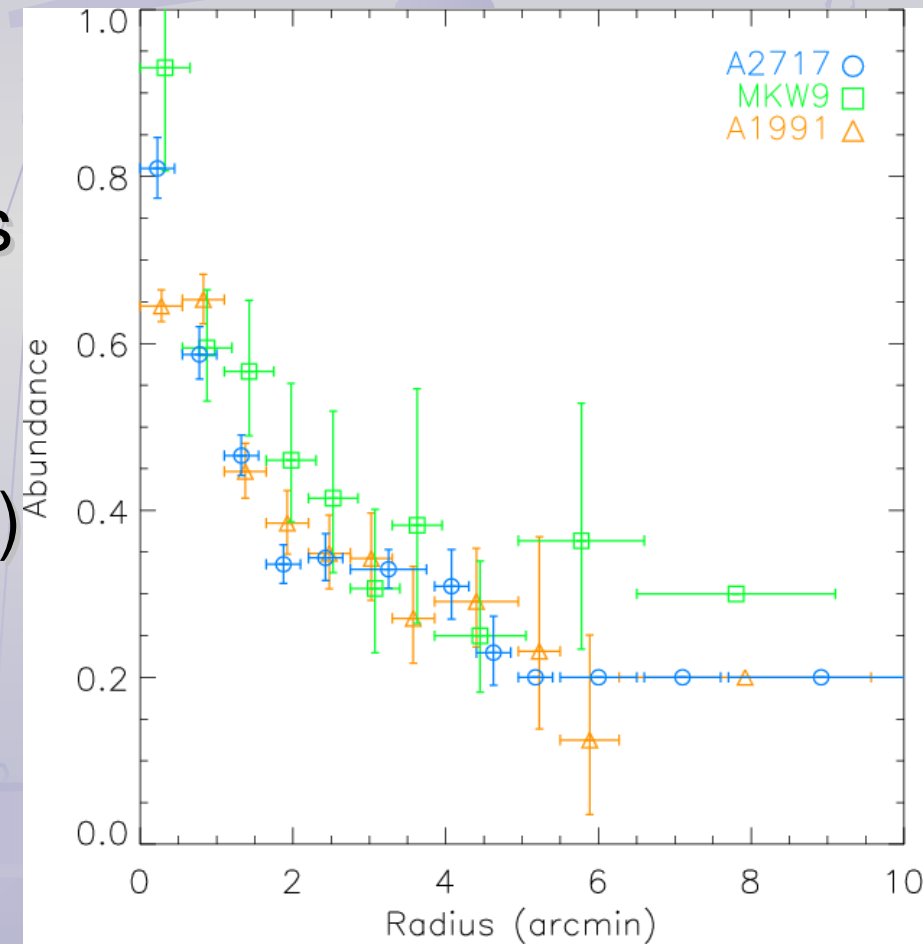


# Projected abundance / radial temperature profiles

- Annular spectral analysis

- Spectra of circular annuli
- Fitted with parameters of temperature ( $\rightarrow$ ), abundance (rel to '89 Anders and Grevesse)

- Projected abundance profiles  $\rightarrow$





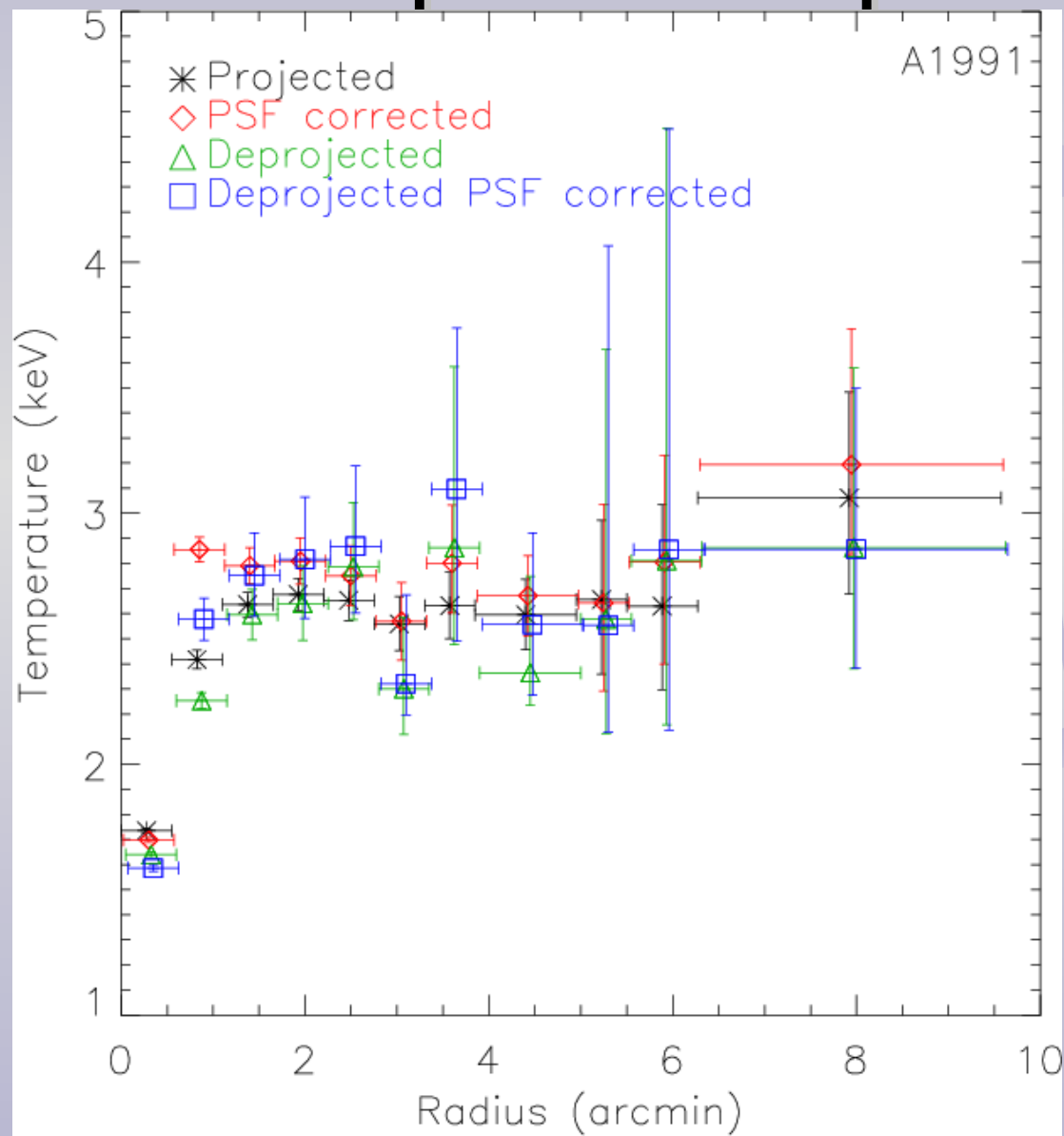
# Radial temperature profiles

- Correction for projection, PSF effects
- Unreal cooling in center of clusters
- Use modeling method, annular spectra modeled with a linear combination of absorbed isothermal MEKAL models –

$$S_i^O(E) = \text{WABS}(N_H^i) \sum_{j=1}^n a_{i,j} \text{MEKAL}(T_j, Z_j).$$

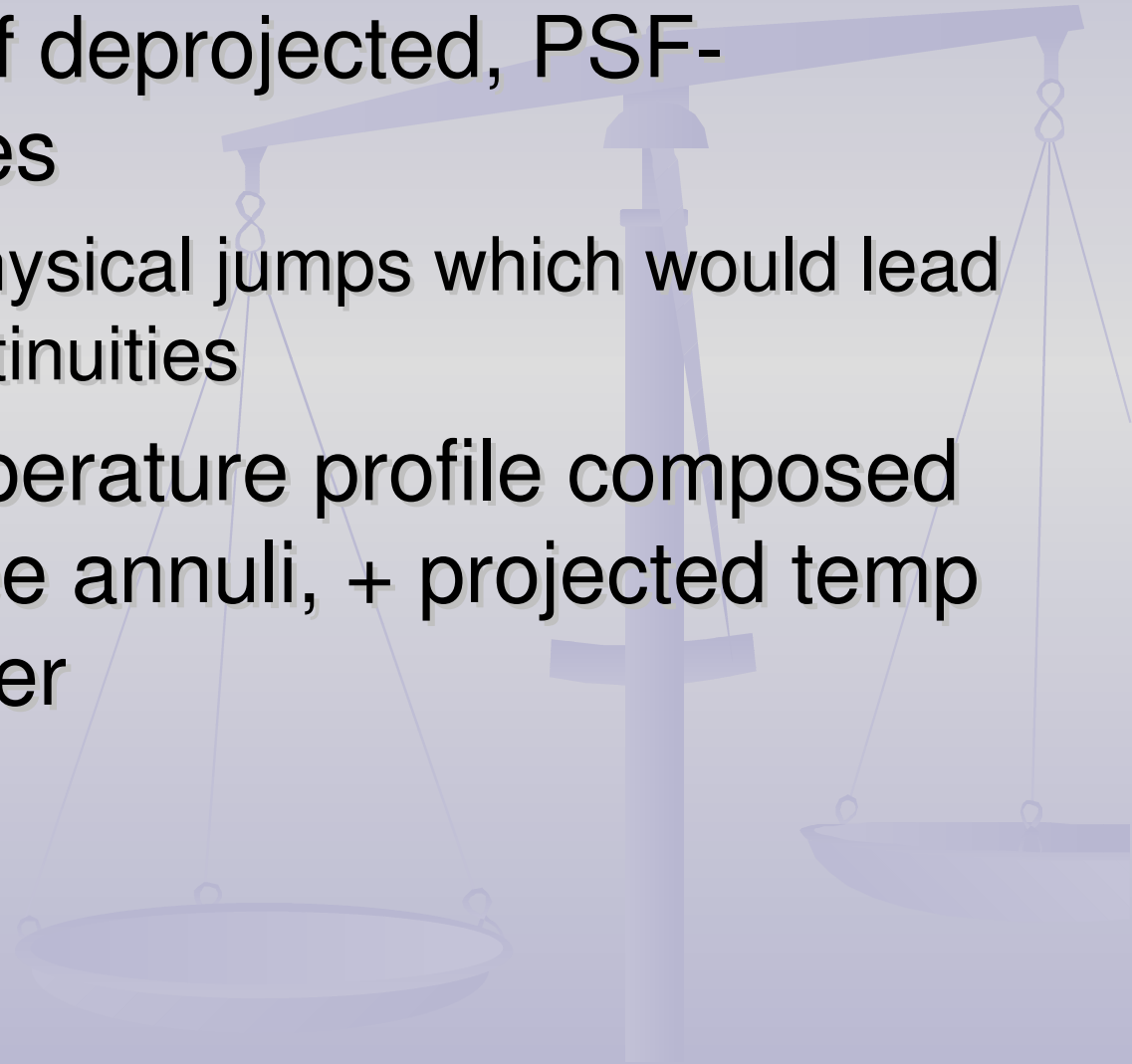
- (PSF) / projection correcting – the  $a_{i,j}$  redistribution coefficients are the EM contributions of (ring j to ring i) / shell j to ring i.

# Radial temperature profiles



# Radial temperature profiles

- Outer regions of deprojected, PSF-corrected profiles
  - Subject to unphysical jumps which would lead to mass discontinuities
- Now use a temperature profile composed of the inner three annuli, + projected temp profiles thereafter



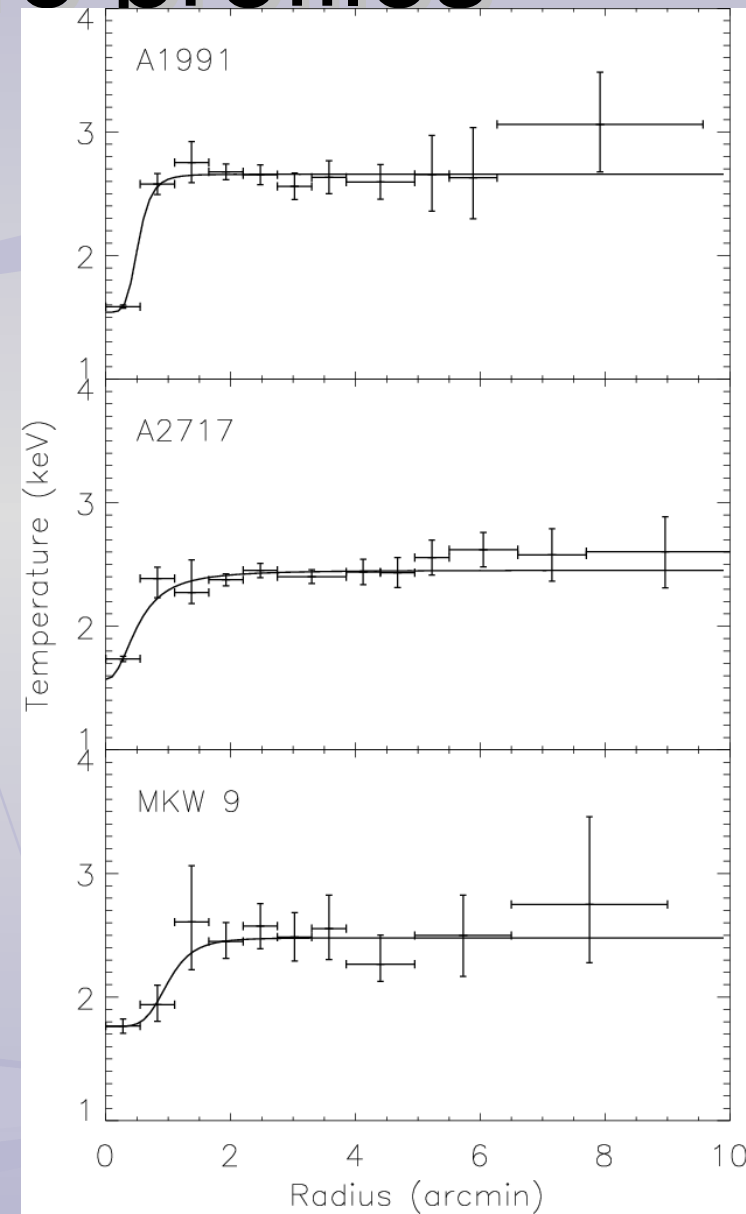
# Radial temperature profiles

- Solid lines are best fit to:

$$T = T_0 + T_1[(r/r_c)^n / (1 + (r/r_c)^n)].$$

- With parameters:

Cluster	T <sub>0</sub>	T <sub>1</sub>	r <sub>c</sub>	n
	(keV)	(keV)		
A1991	1.54	1.12	0 : 52	5.
A2717	1.57	0.88	0 : 52	2.28
MKW9	1.76	0.72	1 : 00	5.



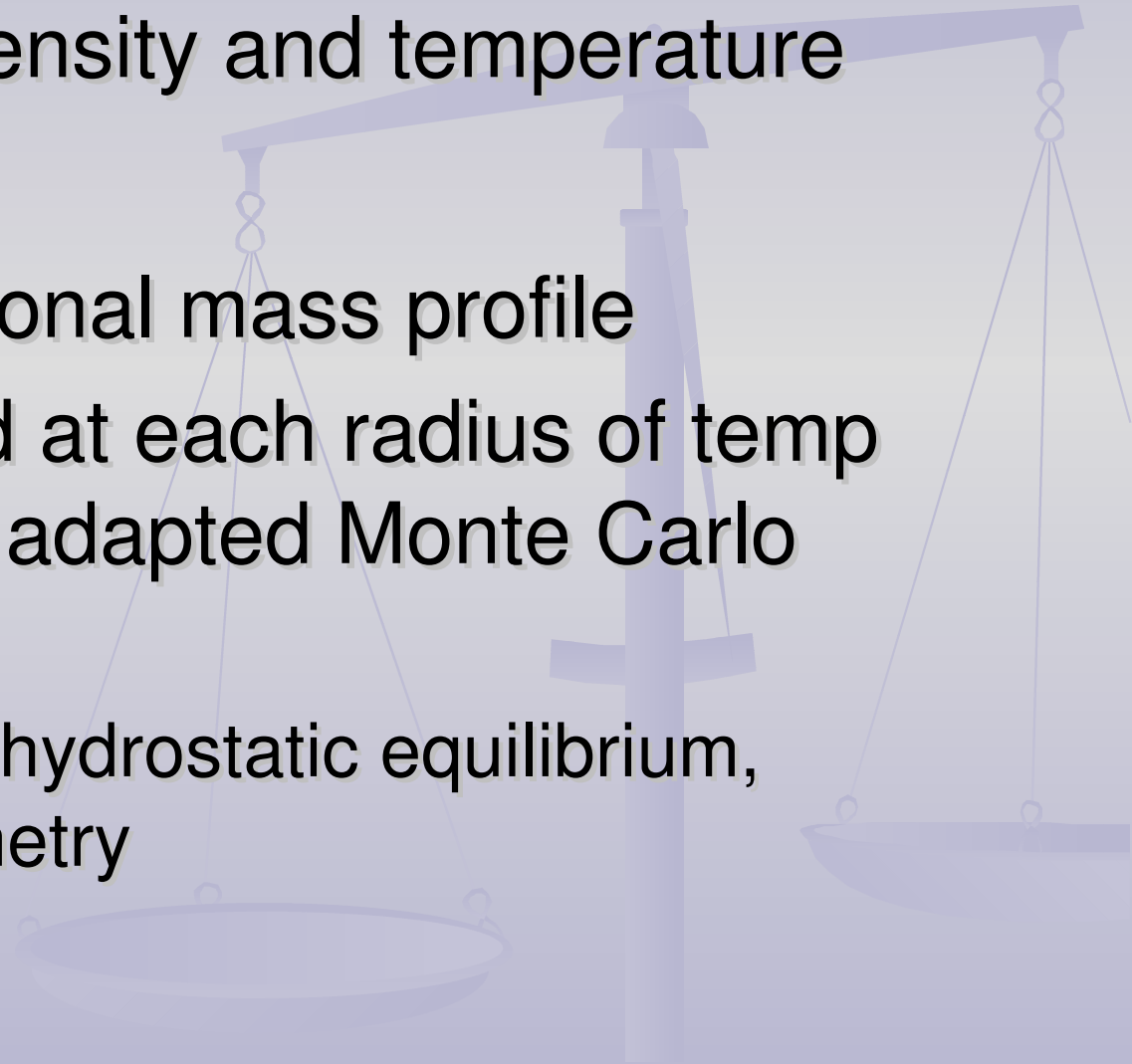
# Global temperature

- Spectra of events in  $0.1 r_{200} \leq r \leq 0.3 r_{200}$ 
  - Higher than  $0.4 r_{200}$ , very little emission
  - Lower than  $0.1 r_{200}$ , cooler gas
- $r_{200}$  from best-fit NFW mass model – later
- Temperature and abundance values:
  - A1991, MKW 9 in agreement
  - A2717 higher than previous, probably better

Cluster	$kT$	$Z$
	(keV)	( $Z_{\odot}$ )
A1991	$2.65^{+0.05}_{-0.05}$	$0.33^{+0.03}_{-0.02}$
A2717	$2.53^{+0.05}_{-0.05}$	$0.34^{+0.02}_{-0.02}$
MKW9	$2.58^{+0.15}_{-0.15}$	$0.37^{+0.07}_{-0.06}$

# Calculating mass profiles

- Combine gas density and temperature profiles
- → total gravitational mass profile
- Mass calculated at each radius of temp profile using an adapted Monte Carlo method
  - Assumptions – hydrostatic equilibrium, spherical symmetry



# Mass profile modeling

- Fitted to density distribution by Navarro et al. (**NFW**)

$$\rho(r) \propto [(r/r_s)(1 + r/r_s)]^{-1}$$

- Parameters: normalization factor, scaling radius  $r_s$ ;

Or, mass  $M_{200}$  and concentration parameter

$$c_{200} = r_{200}/r_s$$

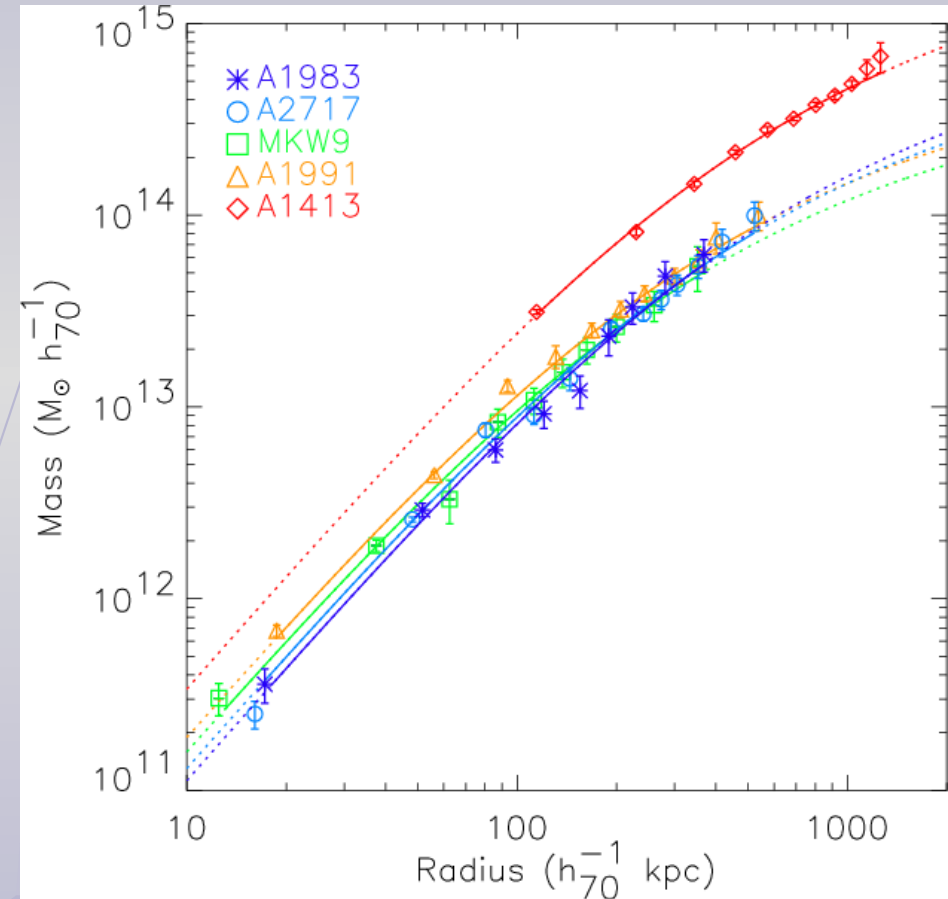
- $M_{200}$ : mass contained within virial radius



# Mass profile modeling.. fits

Parameter	A1991	A2717	MKW9
$\Lambda$ CDMH70			
$C_{200}$	$5.7^{+0.4}_{-0.3}$	$4.2^{+0.3}_{-0.3}$	$5.4^{+0.7}_{-0.7}$
$r_s$ (kpc)	$191^{+19}_{-17}$	$261^{+27}_{-24}$	$186^{+43}_{-34}$
$r_{200}$ (kpc)	1105	1096	1006
$M_{200}$ ( $10^{14} M_{\odot}$ )	1.63	1.57	1.20
$\chi^2/\nu$	9.98/9	15.8/10	4.0/8
SCDMH50			
$c$	$5.6^{+0.4}_{-0.3}$	$4.1^{+0.3}_{-0.2}$	$5.3^{+0.7}_{-0.7}$
$r_s$ (kpc)	$260^{+26}_{-23}$	$358^{+37}_{-33}$	$255^{+61}_{-46}$
$r_{200}$ (kpc)	1466	1466	1358
$M_{200}$ ( $10^{14} M_{\odot}$ )	2.17	2.12	1.63
$\chi^2/\nu$	9.98/9	15.8/10	4.0/8

Results from the NFW fits to the mass profiles



Integrated total gravitating mass profiles,  
1  $\sigma$  errors

# NFW mass profiles & cluster dynamical states

- A2717 not a great NFW fit
  - Maybe halo unrelaxed
- MKW9
  - Unrelaxed.
- Also: use relation between  $M_{200}$  and dark matter velocity dispersion

$$\sigma_{\text{DM}} = 1075 [h(z) M_{200} / (10^{15} h_{100}^{-1} M_{\odot})]^{1/3} \text{ km s}^{-1},$$

and compare with optically-derived velocity dispersions

# NFW mass profiles & cluster dynamical states

- Agreement good
- Esp. for A2717
- So total mass estimates from NFW fits are trustworthy.
- And no cluster is very far from equilibrium.
- Differences in dynamical states don't seem to affect the NFW fit values

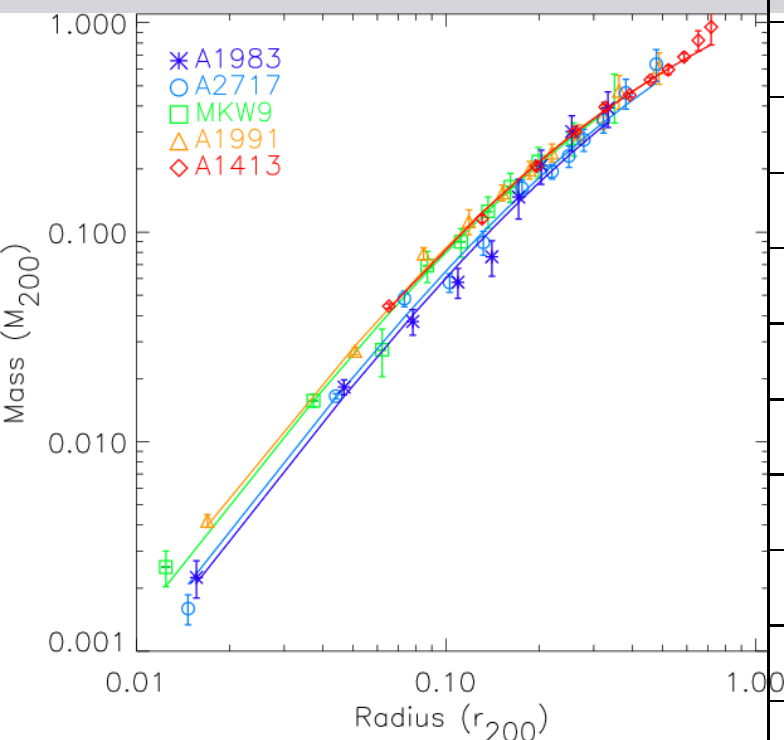
Cluster	$\sigma_{DM}$ km/s	$\sigma_{opt}$ km/s	Reference
A1991	526	$631^{+147}_{-137}$	Girardi et al. (1997)
A2717	520	$541^{+65}_{-41}$	Girardi et al. (1997)
MKW9	474	$579^{+331}_{-337}$	Beers et al. (1995)

Dark matter velocity dispersions calculated from the  $\sigma_{DM} - M_{200}$  relation vs. the optically-derived galaxy velocity dispersions

# Scaled mass profiles

Relative dispersion

Mass profiles



Radius	$\Lambda$ CDMH70		SCDMH50	
	$\sigma / m$	$S_{E1}/C_{E1}$	$\sigma / m$	$S_{E1}/C_{E1}$
Scaled mass: NFW best fit model				
$0.05 r_{200}$	0.18	0.19	0.18	0.19
$0.1 r_{200}$	0.15	0.16	0.15	0.16
$0.3 r_{200}$	0.08	0.09	0.08	0.08
$0.5 r_{200}$	0.04	0.05	0.04	0.04
Scaled mass: interpolated data				
$0.05 r_{200}$	0.18	0.20	0.17	0.18
$0.1 r_{200}$	0.25	0.25	0.24	0.24
$0.3 r_{200}$	0.06	0.02	0.06	0.04
$0.5 r_{200}$	0.13	0.13	0.13	0.12

# Entropy profiles



S vs. T

S vs. M

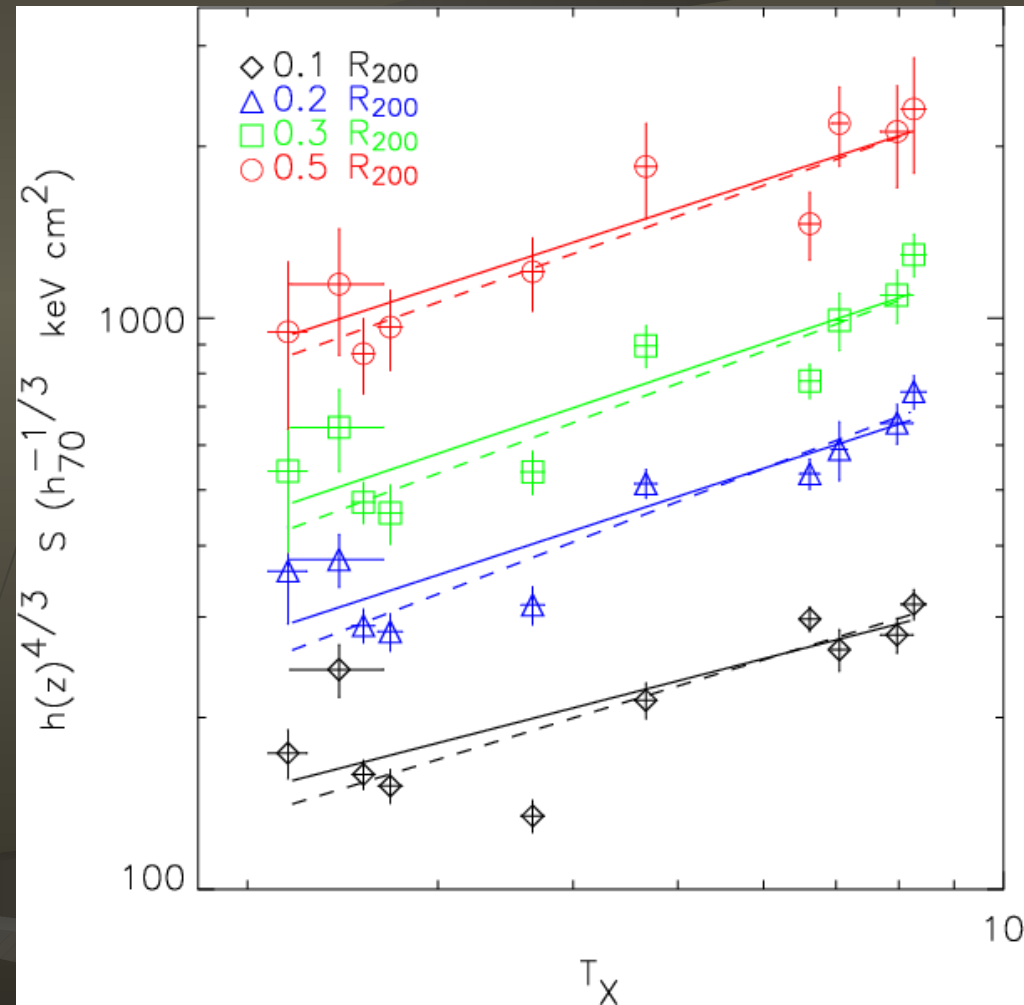
S vs. r, scaled with S-T, S-M

# Entropy – temperature relation

- Self-similar expectation:  $S \propto h(z)^{-4/3} T$
- Data fitted with power law

$$h(z)^{4/3} S_x = A [T_{\text{spec}} / 5 \text{ keV}]^a$$

using three  
regression  
methods



# S vs. T

- At 0.3 R<sub>200</sub>:
  - Gas density / T measurements well constrained – least affected by PSF, projection effect correction problems
  - Outside cooling cores
  - Known from *Chandra* no significant T gradients

$$S_{0.3} \propto T^{0.64 \pm 0.11}$$

- Slope stable after 0.2 R<sub>200</sub>
- Intrinsic scatter largest at 0.1 R<sub>200</sub>

Radius		$\sigma_{\log}$		
$R_{200}$	$\alpha$	raw	stat	int
WLS				
0.1	$0.58 \pm 0.05$	0.079	0.030	0.073
0.2	$0.73 \pm 0.06$	0.058	0.035	0.047
0.3	$0.71 \pm 0.07$	0.074	0.043	0.060
0.5	$0.68 \pm 0.12$	0.070	0.078	-
BCES				
0.1	$0.49 \pm 0.15$	0.082	0.030	0.076
0.2	$0.62 \pm 0.11$	0.063	0.034	0.052
0.3	$0.64 \pm 0.11$	0.078	0.043	0.065
0.5	$0.62 \pm 0.08$	0.074	0.078	-
WLSS				
0.1	$0.47 \pm 0.14$	0.083	0.030	0.077
0.2	$0.67 \pm 0.10$	0.059	0.035	0.048
0.3	$0.69 \pm 0.12$	0.075	0.043	0.061
0.5	$0.68 \pm 0.12$	0.070	0.078	-

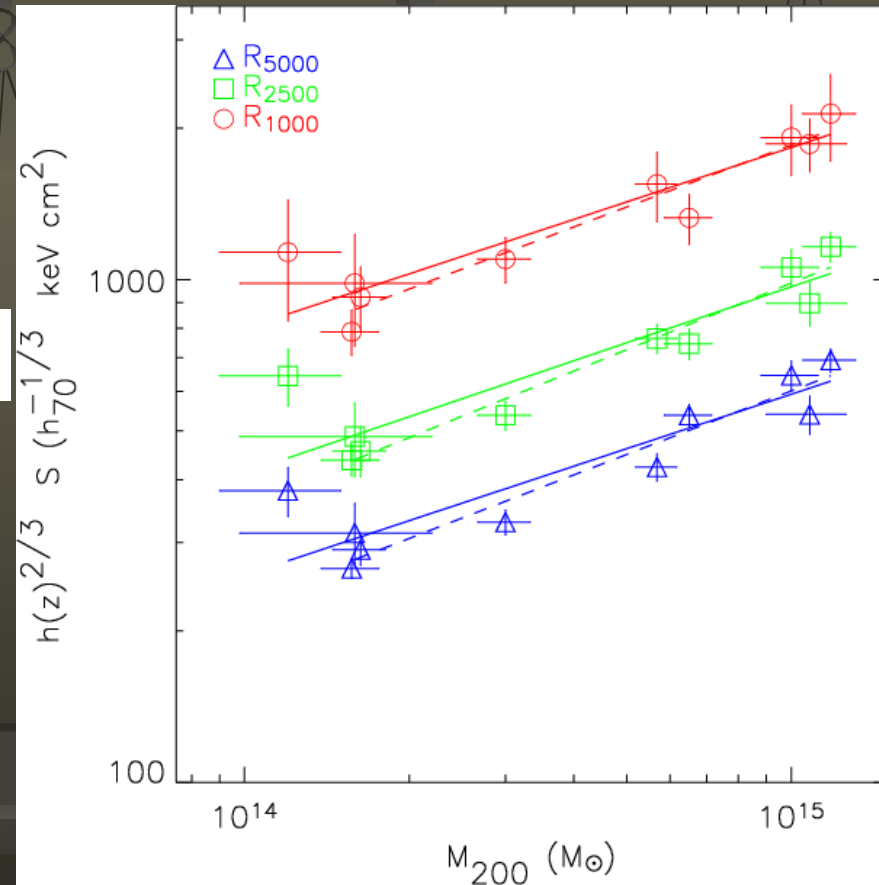
# S – M relation

$$S \propto h(z)^{-2/3} M^{2/3}$$

- Self-similar prediction:
- Best fit is shallower
- Consistent with S – T, M – T relations

$$h(z)^{2/3} S_{\delta} = B_{\delta} \times (M_{200}/5.3 \times 10^{14} M_{\odot})^{\beta}$$

- $M_{200}$  the total mass obtained from NFW fits to the mass profiles





# S vs. M

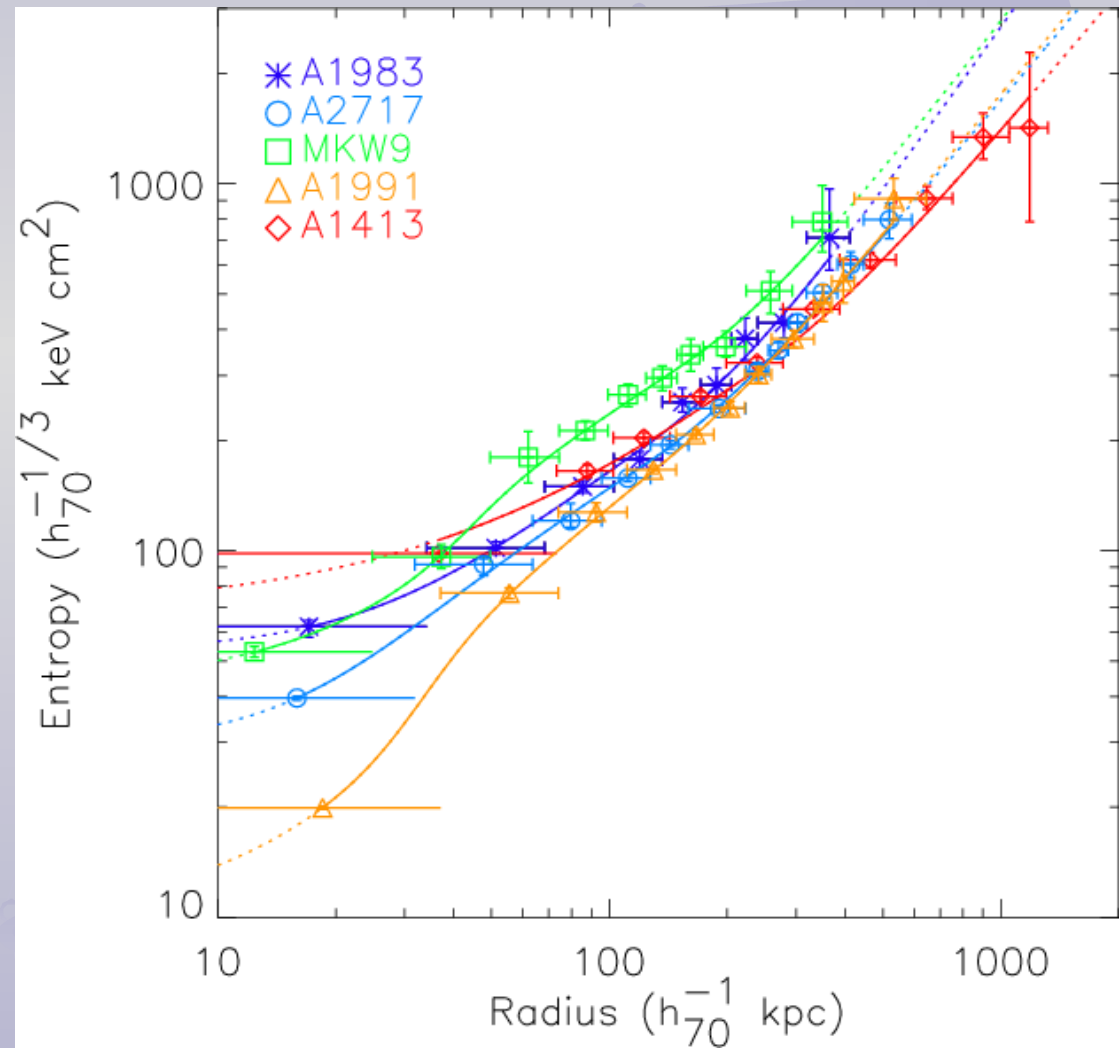
- $S_{2500} \propto h(z)^{-2/3} M_{15}^{0.44}$

- Consistent with  
S – T, S – M  
relations

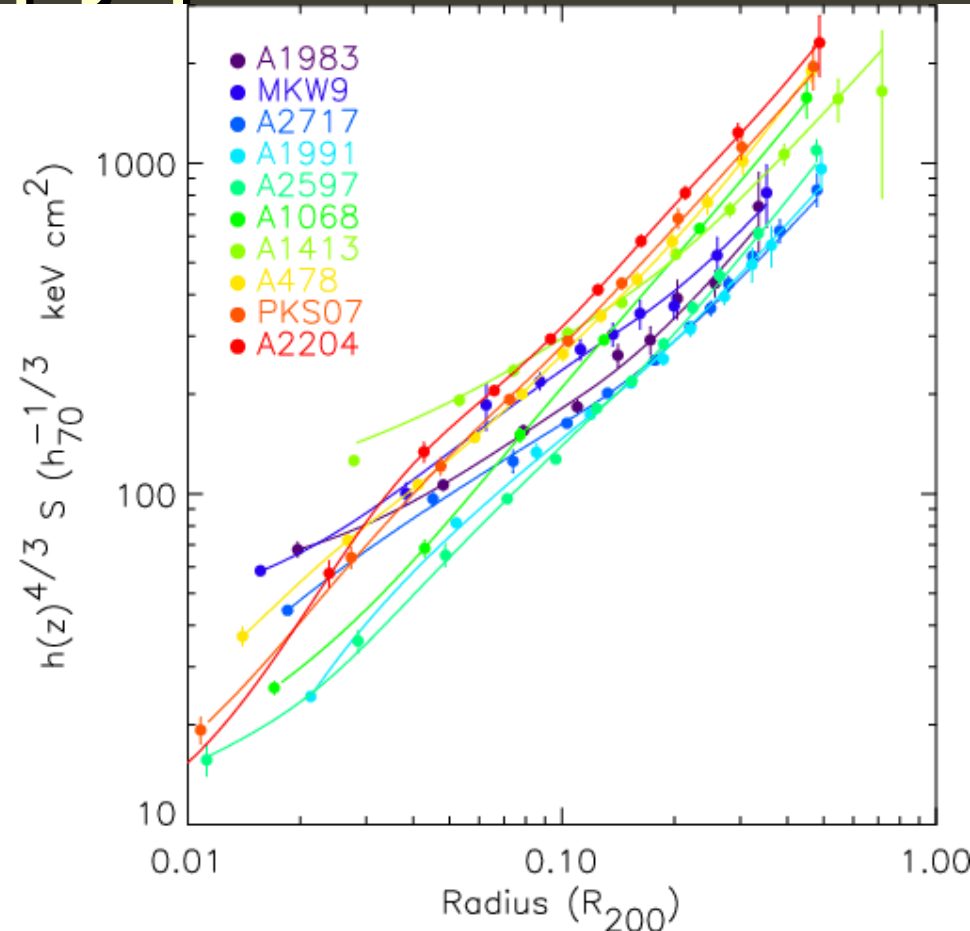
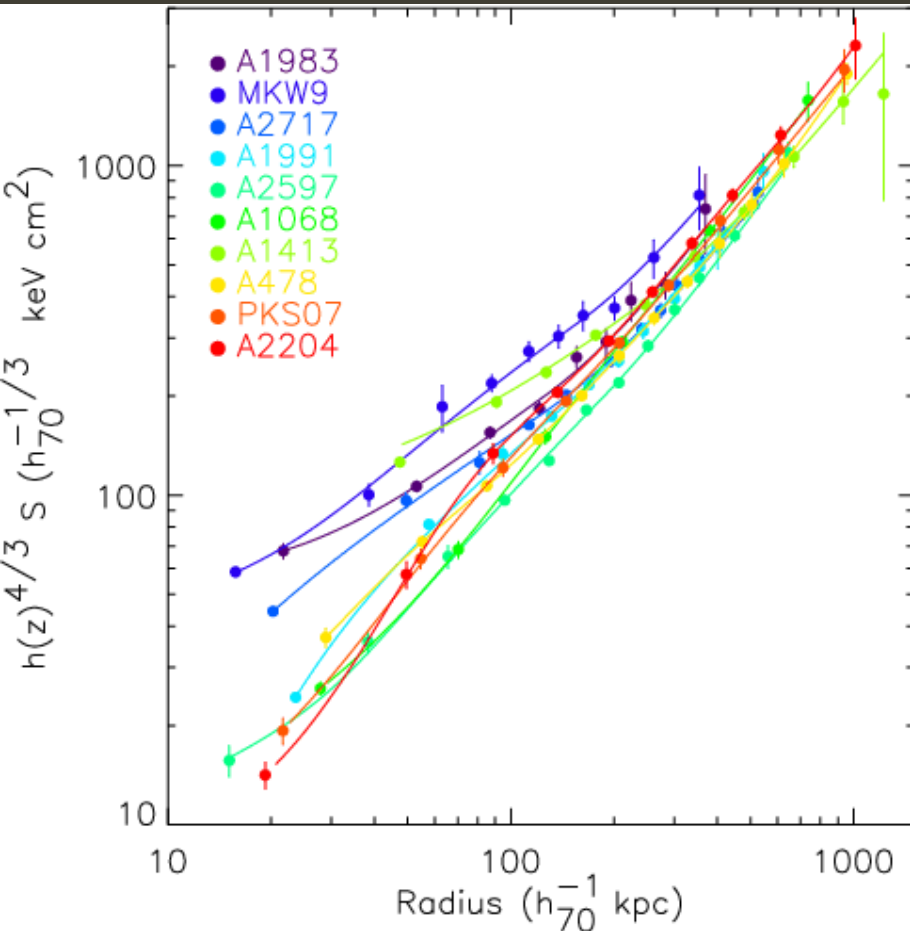
$\delta$	$B$	$\beta$	$\sigma_{\log}$		
			raw	stat	int
	keV cm <sup>-2</sup>				
Full	sample				
5000	471 ± 18	0.36 ± 0.10	0.058	0.034	0.046
2500	765 ± 30	0.37 ± 0.10	0.059	0.041	0.043
1000	1460 ± 47	0.36 ± 0.06	0.059	0.065	-
Excl.	MKW9				
5000	459 ± 37	0.43 ± 0.10	0.039	0.035	0.017
2500	741 ± 43	0.44 ± 0.08	0.035	0.041	-
1000	1430 ± 46	0.41 ± 0.04	0.042	0.063	-

# Raw entropy profiles

- $S = kTn_e^{-2/3}$
- Analytic model gas density profile
- Solid – analytic temp distribution model
- Dotted – observed temp profile



# Raw entropy profiles



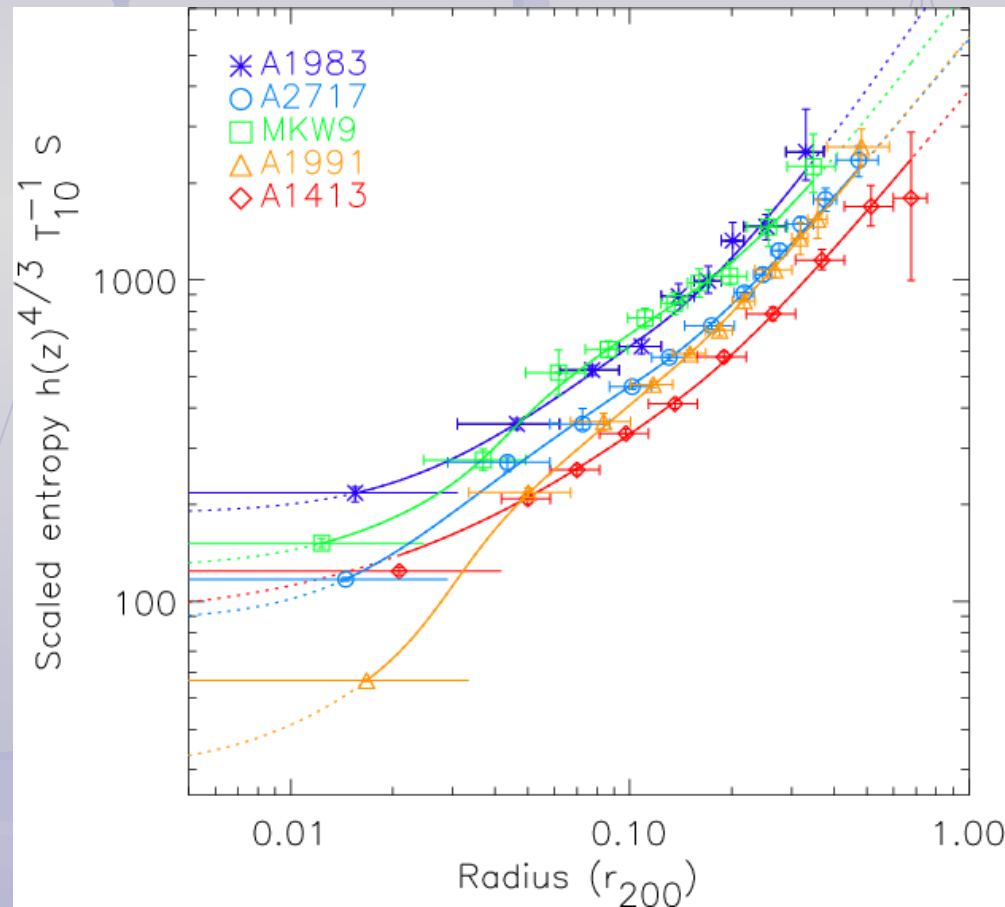
- All profiles increase monotonically with radius
- None have isentropic core

# Scaled entropy profiles

- $T_{10}$  = global temperature in units of 10 keV
- Self-similar scenario:  
clusters form at constant  
density contrast, gas  
follows dark matter

$$\overline{n_e} \propto \overline{\rho_{DM}} \propto \rho_c(z) \propto h^2(z)$$

- Then  $S \propto h(z)^{-4/3} T$  and  
the scaled entropy  
profiles of all clusters  
should coincide

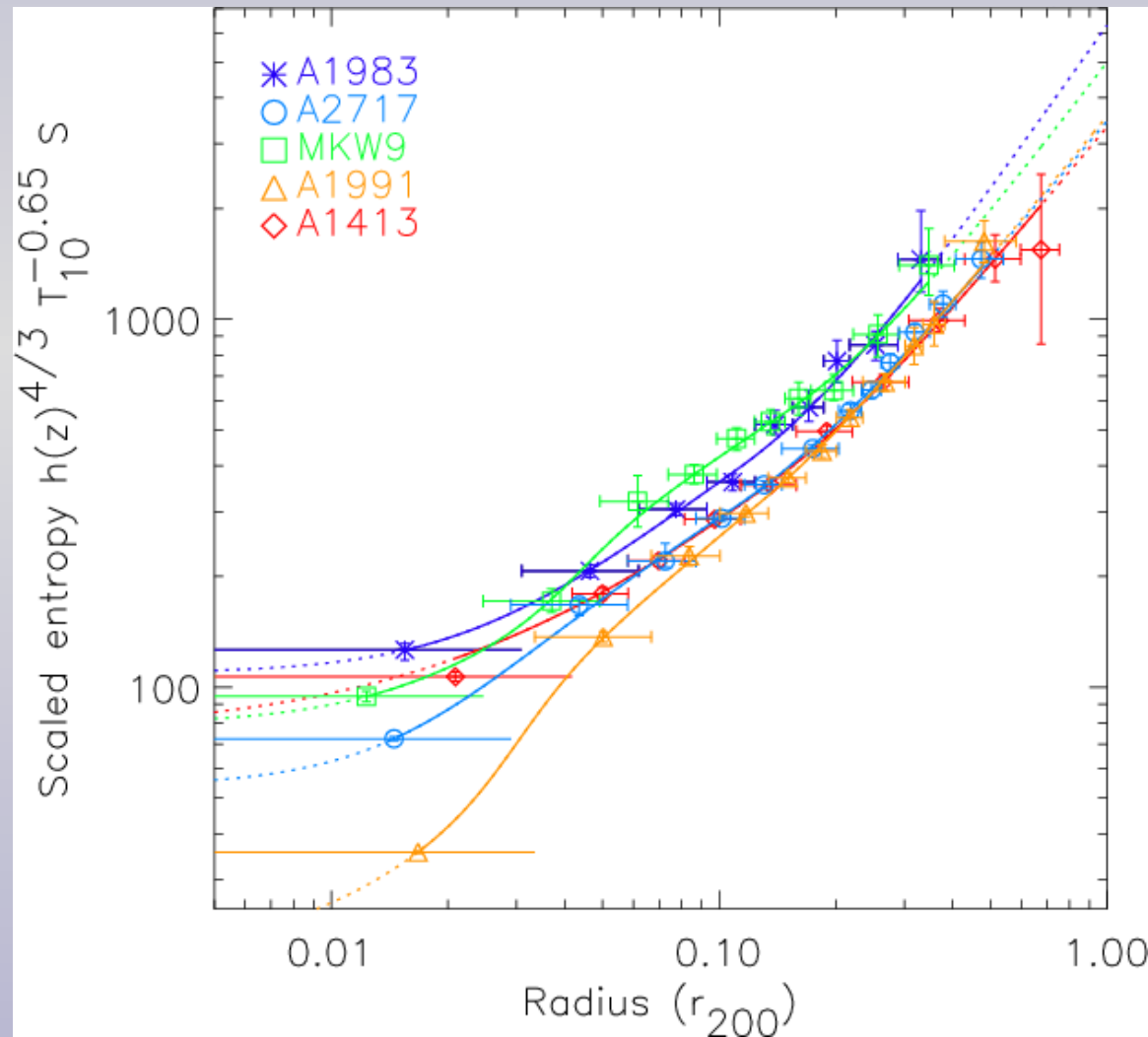


# Scaled entropy profiles

Radius	$\Lambda$ CDMH70		SCDMH50	
	$\sigma$ /m	$S_{\text{BI}}/C_{\text{BI}}$	$\sigma$ /m	$S_{\text{BI}}/C_{\text{BI}}$
Scaled Entropy: $T^{-1}$ scaling				
$0.05 r_{200}$	0.30	0.28	0.28	0.26
$0.1 r_{200}$	0.30	0.29	0.29	0.28
$0.3 r_{200}$	0.30	0.29	0.28	0.26
$0.5 r_{200}$	0.34	0.33	0.31	0.31
Scaled Entropy: $T^{-0.65}$ scaling				
$0.05 r_{200}$	0.22	0.22	0.21	0.20
$0.1 r_{200}$	0.22	0.24	0.21	0.21
$0.3 r_{200}$	0.20	0.26	0.18	0.20
$0.5 r_{200}$	0.24	0.29	0.21	0.19

# Scaled entropy profiles

Better as  $S \propto T^{0.65}$



# Scaled entropy profiles

- Region  $0.05 - 0.1 r_{200}$  well approximated by power law

$$h(z)^{4/3} T_{10}^{-0.65} S(r) = 470 \left( \frac{r}{0.1 r_{200}} \right)^{0.94 \pm 0.14} h_{70}^{-1/3} \text{ keV cm}^2,$$

- Slope close to but shallower than the

$$S \propto r^{1.1}$$

expected from analytical models of shock heating in spherical collapse

# Scaled entropy profiles

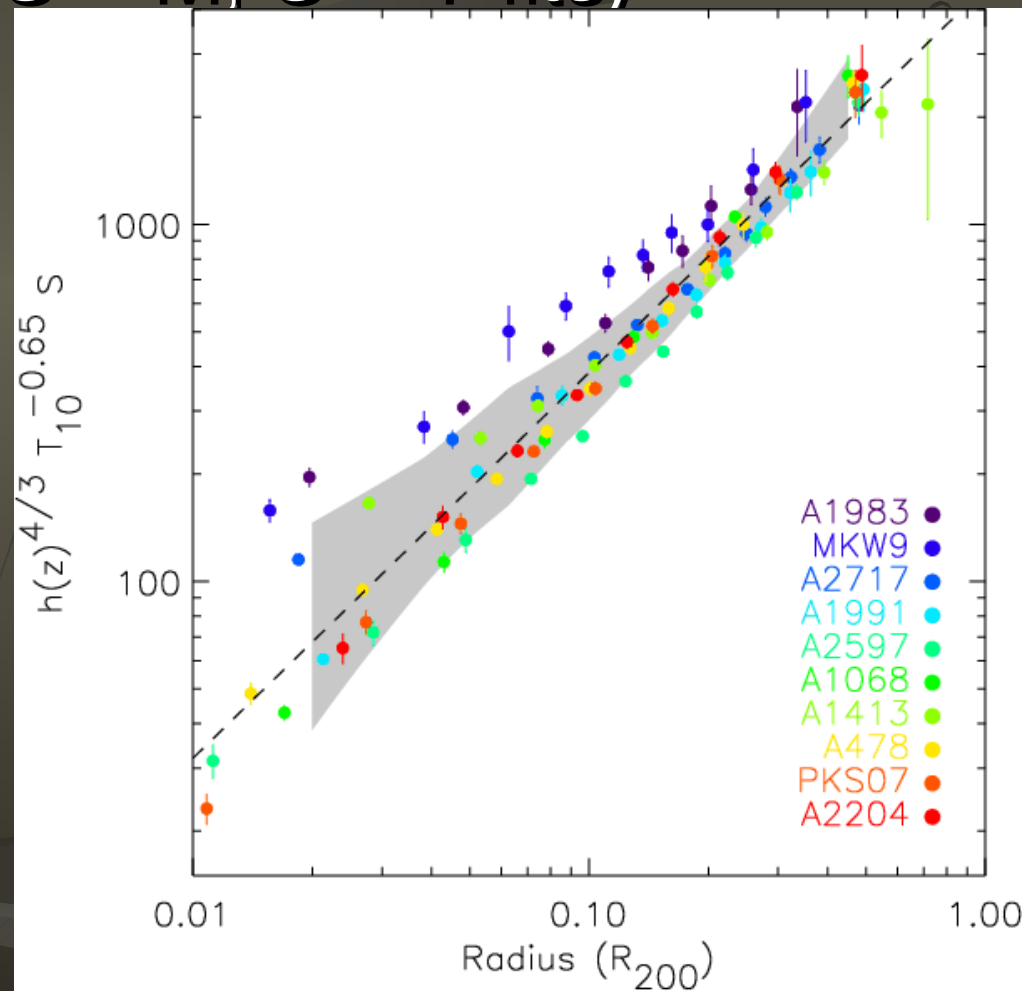
- (scaled by the best  $S - M$ ,  $S - T$  fits)

- $S \propto r^{1.08 \pm 0.04}$

- $S \propto r^{1.1}$

expected for shock  
heating in spherical  
collapse

- Consistent



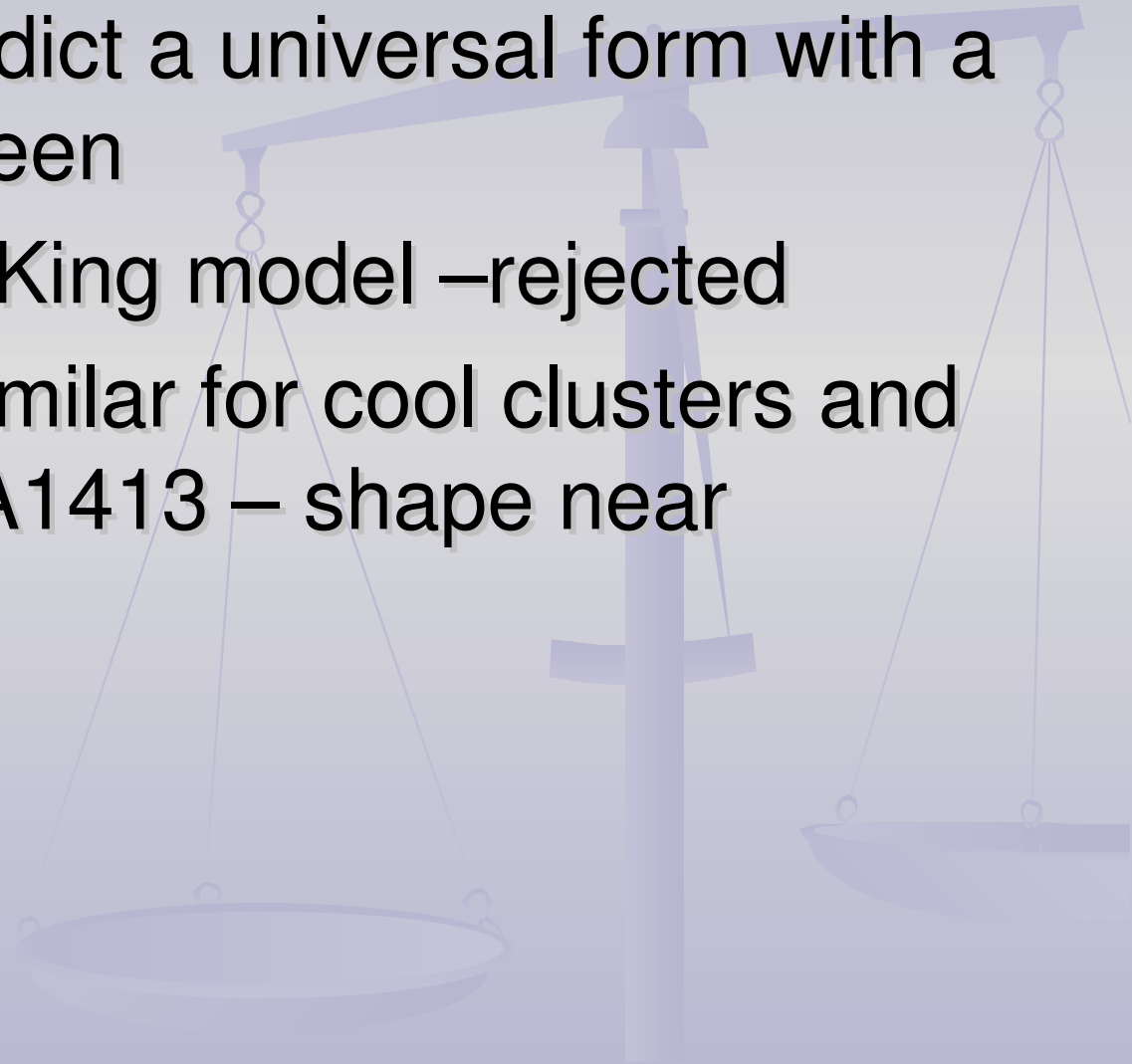


# Discussion



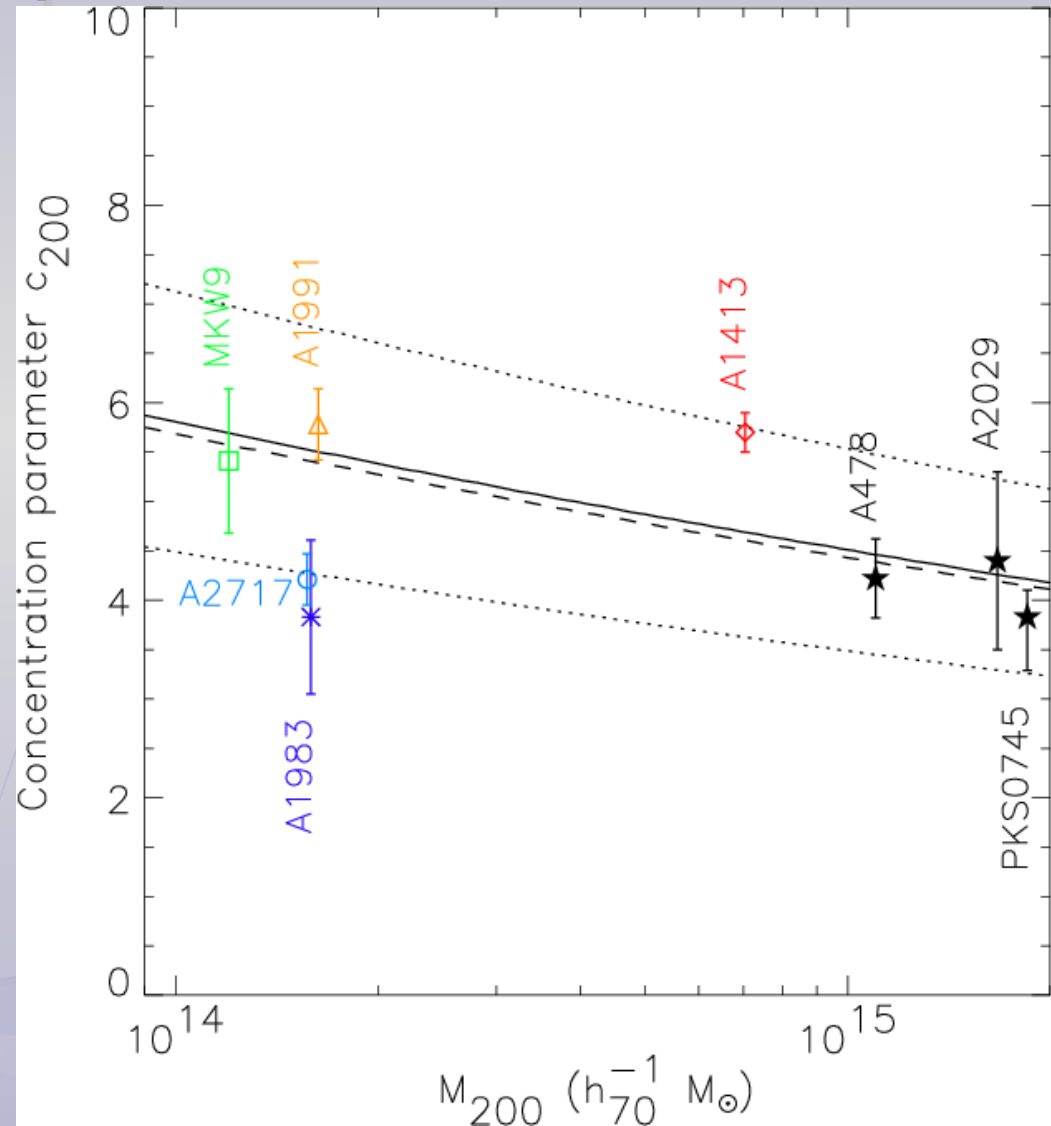
# The gravitational collapse of the dark matter – qualitative check

- Simulations predict a universal form with a central cusp - seen
- NFW – best fit; King model –rejected
- Mass profiles similar for cool clusters and the hot cluster A1413 – shape near universal



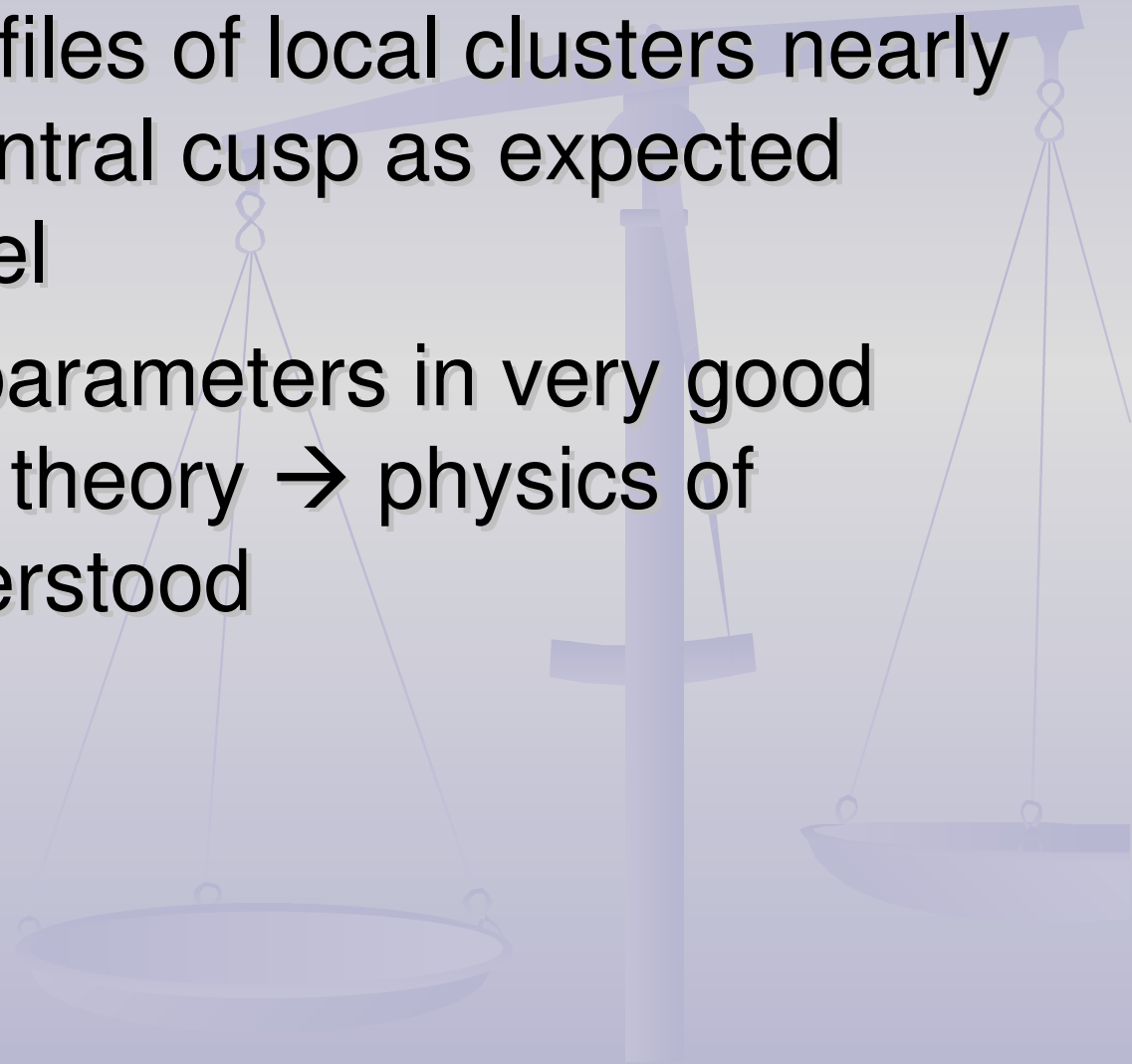
# The gravitational collapse of the dark matter – quantitative check

- Concentration parameters  $c_{200}$  ( $= r_{200}/r_s$ ) should increase for lower mass systems
- Clusters + literature
- Solid line:  $z = 0.0$
- Dashed:  $z = 0.15$



# Dark matter collapse

- Dark matter profiles of local clusters nearly universal, w/ central cusp as expected from NFW model
- Concentration parameters in very good agreement with theory → physics of collapse is understood



# Entropy profiles v. theory

- Departures from self-similar picture – non-gravitational processes
  - Or, due to a flaw in the gravitational collapse model?
  - This study says no.
- Pure cooling / simple preheating models insufficient
  - Spherical preheating – predict a break in  $S - T$  relation and large isentropic cores – rule out
  - Pure cooling models – predict overcooling at odds with observed mass fraction of the stellar component – rule out

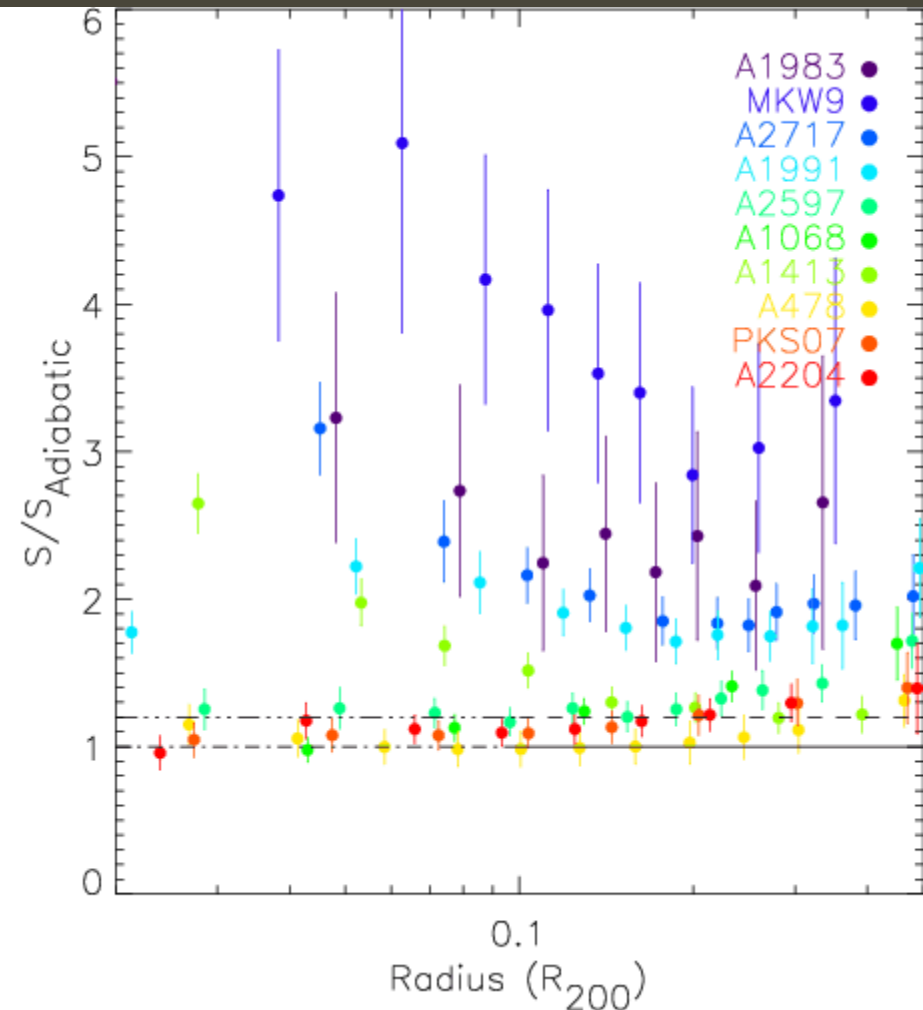
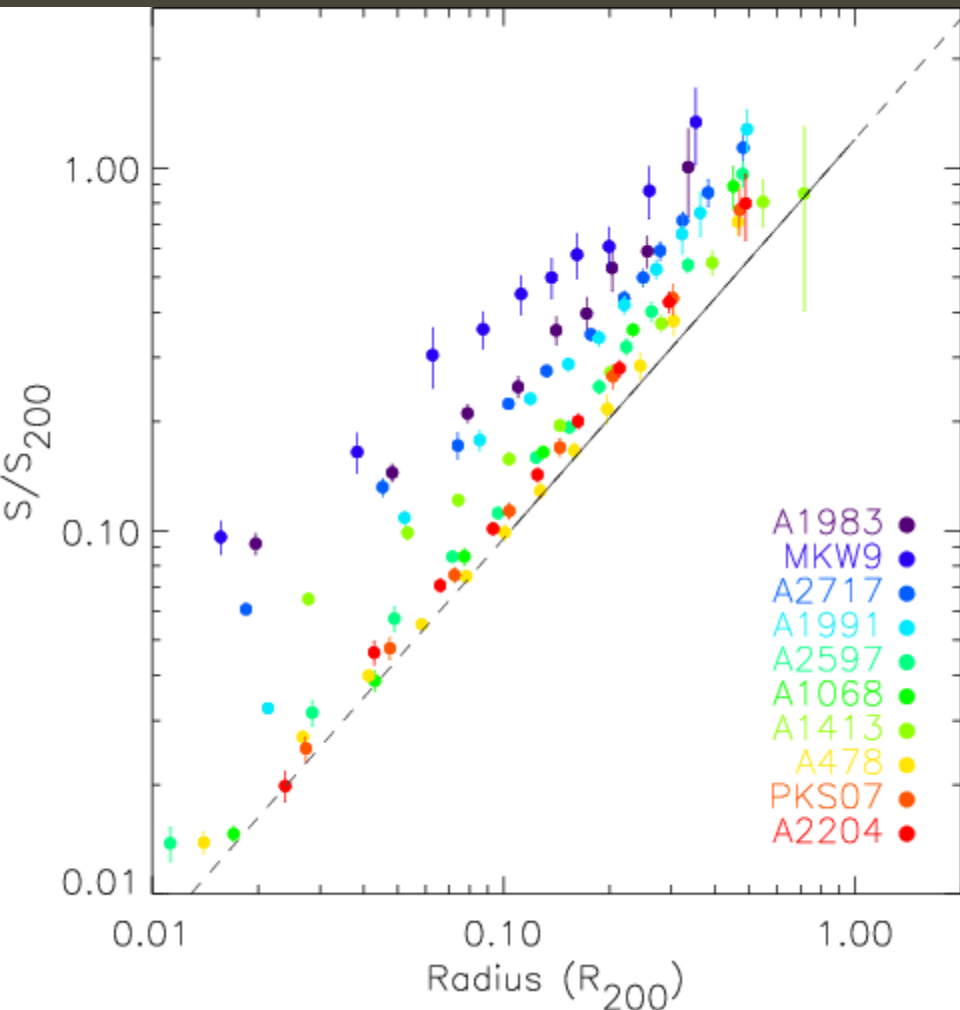
# Entropy profiles v. theory

- Beyond core region ( $r > 0.1 R_{200}$ ):  
profiles self-similar; shape consistent with model but with shallower temperature / mass scaling than expected
- In core region:  
Break of similarity – dispersion increases with decreasing radius

# Modified scaling

- Modified scaling: excess of entropy in low mass objects relative to more massive systems, as compared to the expectation from pure shock heating
- Quantify absolute value of excess, see if an excess is also present for more massive systems
- Adiabatic numerical simulations – Voit, '05

# Modified scaling





# Modified scaling

- Richer systems:
  - Entropy in good agreement with pure gravitational collapse prediction
  - Only ~20% higher
  - Can be accounted for by the difference in observed  $M - T$  relation and modeled
- Poorer systems:
  - $S \sim 2.5$  times higher at  $0.2 R_{200}$  for A1983 than gravitational heating prediction
  - Excess – density of ICM is affected at lower mass
- Entropy boosted at accretion shock

# Modified scaling

- ICM entropy highly sensitive to the density of the incoming gas
- A smoothing of the gas density by preheating in filaments and/or infalling groups would boost entropy production at the accretion shock
- Affects low-mass systems more – accrete smaller halos more affected by smoothing due to preheating
- No isentropic core because the amount of initial preheating is substantially less than the characteristic entropy of the final halo
- Result in self-similarity down to low mass, with modified scaling

# Similarity break in core

- Entropy dispersion  $\sim 60\%$  at  $0.02 R_{200}$
- Six clusters – strong radiative cooling
  - Very self-similar power law profiles, dispersion  $\sim 13\%$  between  $0.01$  and  $0.1 R_{200}$
  - Consistent with quasi-steady-state models that include radiative cooling
- Four clusters – shallower entropy profiles
  - AGN energy input
  - Strong bursts / *weak shocks*
  - Old merging events – mixing high / low S gas

# Conclusion

- Confirmed physics of dark matter collapse is understood
- Entropy profiles – in  $0.05 r_{200} \leq r \leq 0.5 r_{200}$ , self-similar but scale with  $T$ ,  $r$  shallower than gravity-only, pure shock heating model

$$S(r) \propto T^{0.63 \pm 0.05} (r/r_{200})^{0.94 \pm 0.14}$$

- Large dispersion in  $r \leq 0.05 r_{200}$  – variety of cooling core histories

# Conclusion

- Entropy scales with temperature:

$$S_{0.3} \propto T^{0.64 \pm 0.11}$$

- With radius:

$$S \propto r^{1.08 \pm 0.04}$$

- Modified scaling thought due to smoothing of accreted gas density by preheating
  - Would affect low-mass systems more, as seen
- Large dispersion in core thought due to
  - Some clusters are cooling flow clusters
  - Some – energy coming from weak shocks from AGN activity, effects of old mergers

- The end

